

**<sub>1</sub> Magnetic field fluctuation features at Swarm's  
<sub>2</sub> altitude: a fractal approach**

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3 The ESA Swarm mission provides a qualitatively new level of observational  
4 geomagnetic data, which allows us to study the spatial features of magnetic  
5 field fluctuations, capturing their essential characteristics and at the same  
6 time establishing a correlation with the dynamics of the systems responsi-  
7 ble for the fluctuations. Our study aims to characterize changes in the scal-  
8 ing properties of the geomagnetic field's spatial fluctuations by evaluating  
9 the local Hurst exponent, and to construct maps of this index at the Swarm's  
10 altitude ( $\sim 460$  km). Since a signal with a larger Hurst exponent is more  
11 regular and less erratic than a signal with a smaller one, the maps permit  
12 us to localize spatial structures characterized by different scaling properties.  
13 This study is an example of the potential of Swarm data to give new insights  
14 into ionosphere-magnetosphere coupling; at the same time, it develops new  
15 applications where changes in statistical parameters can be used as a local  
16 indicator of overall magnetospheric-ionospheric coupling conditions.

## 1. Introduction

17 It is well-known that the magnetic field observed at or near the Earth's surface is not  
18 constant, but affected by variations on different temporal and spatial scales [*Courtillot*  
19 *and Le Mouél*, 1988].

20 The magnetic field's temporal variations with periods from seconds to several hundred  
21 minutes are the result of processes related to the interaction between the solar wind and  
22 the Earth's magnetic field. As a result of this interaction, a considerable amount of  
23 energy is released, giving rise to a number of important phenomena in the magnetosphere  
24 and upper atmosphere. Examples include large-scale plasma motions, electric currents,  
25 aurorae, and disturbances of the neutral and ionized upper atmosphere [*Prölss*, 2006].  
26 Some of these phenomena can be identified readily in magnetograms, recorded both on  
27 the ground at geomagnetic observatories and in near-Earth space by satellites orbiting  
28 our planet.

29 Within this system, the European Space Agency's Swarm mission provides a good op-  
30 portunity for an average study of the magnetic fields of external origin on a global scale.  
31 The multipoint measurements of the Swarm constellation mission provide a qualitatively  
32 new level of observational data, which allows us for the first time to solve the spatio-  
33 temporal features of the magnetic field fluctuations of external origin, to capture their  
34 essential characteristics at the Swarm's altitude ( $\sim 460$  km), and, at the same time, to  
35 attempt to establish a correlation with the dynamics of the system responsible for such  
36 fluctuations. However, some difficulties arise when this type of data is used. The exter-  
37 nal contributions to the geomagnetic field are ordered primarily in a local time frame;

38 satellites in a polar orbit, like those of the Swarm constellation, can obtain a reasonably  
39 dense sampling of the internal components within a few days, but fails to provide ade-  
40 quate spatial coverage of the external contributions, because of the slow orbital precession  
41 through local time. Thus, the Swarm constellation configuration does not allow real-time  
42 monitoring of the magnetic field's spatial fluctuations, but only an average study of them.

43 Recently, the study of magnetic field fluctuations of external origin, recorded both  
44 on the ground and in different parts of the Earth's magnetosphere, has been addressed  
45 by many researchers because it has permitted a better understanding of the complex  
46 magnetospheric dynamics in response to solar wind changes. The analysis of magnetic field  
47 fluctuations and in particular of the geomagnetic indices has played a crucial role in several  
48 works [*Tsurutani et al.*, 1990; *Consolini et al.*, 1996; *Consolini and Chang*, 2001; *Sharma et*  
49 *al.*, 2001; *Uritsky et al.*, 2002; *Consolini et al.*, 2005, 2008] where the nonlinear properties of  
50 the magnetospheric dynamics have been discussed with special attention to the occurrence  
51 of chaos, turbulence, and criticality. However, if the analysis of geomagnetic indices  
52 permits us to study temporal changes of the magnetic field fluctuations, the measurements  
53 of the Swarm constellation provide a unique opportunity to study the spatial features of  
54 magnetic field fluctuations at high resolution. The vector field magnetometer (VFM)  
55 produces measurements of the field's vector at a sampling rate of 1 Hz which, with a  
56 satellite speed of 7.6 km/s, corresponds to a spatial resolution of about 7.6 km along the  
57 orbital track.

58 We use the local Hurst exponent for investigation of the scaling features of geomagnetic  
59 field fluctuations at temporal scales below 40 s (corresponding to spatial scales below  $\sim 300$

60 km) because this quantity, which is a measure of the way in which a data series varies  
61 in time, can be used to obtain significant results for the characterization of dynamical  
62 systems. The local Hurst exponent can be used to characterize the persistence of a  
63 system, e.g., whether the sign of the fluctuations will remain the same (persistent) or  
64 change (anti-persistent) from one point to the next, thus providing information on the  
65 existence of localized spatial structures. In recent years, there has been increasing interest  
66 in the analysis of the Hurst exponent of geomagnetic signals. However, we have found  
67 no studies, which analyze satellite data to reconstruct maps of the Hurst exponent to  
68 characterize the spatial scaling features of geomagnetic field fluctuations.

69 The aim of this letter is to investigate the spatial distribution of the local Hurst exponent  
70 in high-latitude regions for two different geomagnetic activity levels and to attempt an  
71 interpretation in terms of spatial fluctuation structures, and physical processes responsible  
72 for them.

## 2. Dataset Description

73 The present work focuses on the analysis of the fluctuations of the horizontal com-  
74 ponent of the Earth's magnetic field from 1<sup>st</sup> January 2014 to 30<sup>th</sup> June 2014. These  
75 time series have been computed using the calibrated and corrected measurements of the  
76 vector magnetic components in the North-East-Center local Cartesian coordinate frame  
77 recorded by one of the three satellites of the Swarm constellation (Swarm A) (product of  
78 Swarm/VFM and ASM: Swarm level 1b (LR)(MAGA\_LR)). The time interval analysed  
79 contains both periods of relatively low geomagnetic activity and periods characterised by  
80 the occurrence of moderate activity. Being interested in the high-latitude regions (lati-

81 tude higher than  $50^\circ$  N), we choose the Auroral Electrojet (AE) index to discriminate the  
82 different magnetospheric activity levels. The AE index is designed to provide a global,  
83 quantitative measure of auroral zone magnetic activity produced by enhanced ionospheric  
84 currents flowing below and within the auroral oval. In this work we have selected two  
85 different activity levels:  $AE < 60$  nT (quiet geomagnetic activity level) and  $AE > 80$  nT  
86 (disturbed geomagnetic activity level) according to the statistical features of AE-index  
87 [*Consolini and De Michelis, 1998*]. One-minute data of this index have been downloaded  
88 from OMNI data set, where the available data stop at 30<sup>th</sup> June 2014, which justified our  
89 selection of Swarm's dataset.

### 3. Method of Analysis

90 For our investigation of magnetic spatial fluctuation features, we use the Hurst exponent  
91  $H$ , a measure of the persistence features of a time series. The value of the Hurst exponent  
92 lets us ascertain whether the analyzed time series has an anti-persistent or persistent  
93 character. It has been shown that a Hurst exponent value between 0 and 0.5 exists  
94 for signals with *anti-persistent character* of fluctuations. This means that a positive  
95 (negative) fluctuation will tend to be followed by a negative (positive) one. Conversely, a  
96 Hurst exponent value between 0.5 and 1 indicates a *persistent character* of fluctuations,  
97 so that a positive (negative) fluctuation will tend to be followed by another positive  
98 (negative) one – that is, the signal is trending. The larger the  $H$  value is, the stronger the  
99 trend. In other words, the underlying spatial structure is governed by a positive feedback  
100 mechanism of the fluctuations. Lastly, a Hurst exponent value equal to 0.5 indicates  
101 that there is no correlation between the repeated increments. This value also marks the

102 transition between anti-persistent and persistent behavior in a signal. Thus, physically  
103 speaking, investigation of the Hurst exponent may allow detection of localized spatially  
104 coherent structures.

105 Since the geomagnetic field fluctuations do not exhibit a simple global scaling behavior  
106 which can be described using a single scaling exponent, it is necessary to introduce a  
107 multitude of scaling exponents [*Consolini et al.*, 1996; *Consolini and De Michelis*, 1998;  
108 *Wanliss*, 2005]. In many physical systems the scaling features acquire a local character,  
109 which can have a dependence on the amplitude of fluctuations for multifractal objects,  
110 or a time dependence for multifractional signals. Here, instead of using the standard  
111 global Hurst exponent to characterize the properties of these time series, it is better  
112 to introduce a local Hurst exponent because the scaling properties of time series under  
113 investigations cannot be considered constant. Indeed, it is of extreme importance to  
114 correctly quantify the long-range correlations of the geomagnetic time series in order to  
115 gain a deep understanding of the complex system dynamics that gives rise to the recorded  
116 geomagnetic signal.

117 In the past many different techniques and methods, based on the analysis of time series  
118 features in the real or Fourier space, have been proposed to estimate the local Hurst  
119 exponent of a time series [*Holschneider*, 1988; *Bacry et al.*, 1993; *Muzy et al.*, 1994; *Peng*  
120 *et al.*, 1992, 1995; *Abry et al.*, 2000; *Alessio et al.*, 2002]. In this paper, we employ an  
121 alternative method based on the detrended 1<sup>st</sup>-order structure function  $S_1(\tau)$ , which for  
122 a signal  $x(t)$  defined over an interval  $T$  is given by

$$S_1(\tau) = \langle |x(t + \tau) - x(t)| \rangle_T, \quad (1)$$

123 where  $\tau$  is a time separation, and  $\langle \dots \rangle_T$  indicates time averaging over the interval  $T$ .  
 124 This 1<sup>st</sup>-order structure function exhibits a power law behavior as a function of the time  
 125 separation  $\tau$  when we deal with a scale invariant signal  $x(t)$ , so that  $S_1(\tau) \sim \tau^H$  where  $H$   
 126 is the Hurst exponent. Thus, the analysis of the scaling features of the 1<sup>st</sup>-order structure  
 127 function exponent provides an efficient method for characterizing the correlative structure  
 128 of a signal as an empirical approximation to the Hurst exponent.

129 The method used in our analysis can be summarized as follows. Given a time series  
 130  $y(t)$ , we consider a time interval  $[t_0 - T/2, t_0 + T/2]$ , where  $T$  is at least 10 times larger  
 131 than the maximum scale  $\tau$  which we want to investigate. In the selected time interval,  
 132 we detrend the time series by computing the average long-term trend using a 7<sup>th</sup>-order  
 133 polynomial fit  $p(t)$ . In this way, we can construct a new detrended time series  $x(t)$ ,

$$x(t) = y(t) - p(t), \quad \forall t \in [t_0 - T/2, t_0 + T/2], \quad (2)$$

134 to which we apply the structure function analysis. We call the complete procedure as  
 135 *detrended structure function analysis* (DSFA).

136 In the case of Swarm's data, we have calculated the local Hurst exponent  $H$  for fluctua-  
 137 tions in the range from 1 s to 40 s, by considering a moving window of 400 s (i.e.,  $T = 400$   
 138 s in our case). The choice of a maximum timescale of 40 s is motivated by the require-  
 139 ments to get local spatial information on the magnetic field fluctuations. Indeed, using  
 140 the Swarm's satellite velocity and assuming that the evolution time of spatial structures

141 is longer than the transit time, these temporal scales roughly correspond to investigation  
142 of spatial fluctuations from 7.6 km up to  $\sim 300$  km and relate to fluctuations in the mag-  
143 netohydrodynamic (MHD) domain: the ion-gyroperiod is  $T_\Omega = 1/\Omega_i \ll 1$  s, and both the  
144 ion-inertial length  $\eta_i$  and the ion-Larmor radius  $r_L$  are smaller (much smaller) than 7.6  
145 km for the typical ionospheric plasma parameters at Swarm's altitude ( $\sim 460$  km).

146 Typical relative error in the estimation of the local Hurst exponent is 4% with a 95%  
147 confidence. This error has been evaluated using a Monte Carlo simulation. Figure 1 shows  
148 an example of the behavior of the 1<sup>st</sup>-order structure function  $S_1(\tau)$ . A clear power law  
149 dependence is recovered at the scales investigated here, thus assessing the scale invariance  
150 of the analysed signal in the considered range of timescales.

#### 4. Results and Discussion

151 As described in the previous section, we employ DSFA analysis to determine the sta-  
152 tistical nature of our signal. An example of our results is shown in Figure 2, where the  
153 computed local Hurst exponent values are plotted for one Swarm A orbit and reported  
154 with the corresponding values of AE-index, horizontal component of the geomagnetic field  
155 and satellite's geomagnetic latitude. As shown in Figure 2, the character of the analyzed  
156 time series is the result of a superposition of structures (set of fluctuations) characterized  
157 by different values of the local Hurst exponent in the interval  $[0, 1]$ . Consequently, during  
158 the selected period the analyzed time series are characterized at scales below 40 seconds,  
159 both by fluctuations that tend to induce stability within the system, and by fluctuations  
160 with a persistent behavior implying a system dynamics governed by a positive feedback  
161 mechanism. This sample is chosen to better assess the potential of the local Hurst expo-

162 nent to delineate and characterize transitions in magnetograms due to dynamical changes  
163 of the scaling fluctuation features on spatio-temporal scales.

164 Figure 3 shows polar view maps of the local Hurst exponent values over the polar region  
165 in the northern hemisphere for different geomagnetic activity levels. In the top panel of  
166 Figure 3, the spatial distribution of the local Hurst exponent values is reported for the  
167 quiet geomagnetic activity level ( $AE < 60$  nT), while the bottom panel shows the same  
168 quantity during disturbed intervals ( $AE > 80$  nT). The magnetic field fluctuations with a  
169 persistent character are shown in blue, and those with an anti-persistent character in red.  
170 As can be seen from the image in Figure 3 (on the top) there is a strong asymmetry in  
171 the daily character of the magnetic field fluctuations during the geomagnetic quiet level.  
172 At mid/high geomagnetic latitude between  $50^\circ$  N and  $\sim 70^\circ$  N magnetic field fluctuations  
173 show an anti-persistent character in the sunlit hemisphere ( $03 < MLT < 19$ ) and a  
174 mostly persistent one in the dark sector. This difference in magnetic field fluctuation  
175 character reflects the morphology and dynamics of that part of the ionosphere which is  
176 crossed by the satellite. This region can, to varying degrees, be influenced by solar EUV  
177 radiation, energetic particle precipitation, diffusion, thermospheric winds, electrodynamic  
178 drifts, polar wind escape, and so on, which have a different behavior in the dark and  
179 sunlit hemisphere. Moreover, the high-latitude ionosphere differs significantly from its  
180 mid-latitude counterpart because, given the geometry of the geomagnetic field, which is  
181 to a first approximation dipolar, the high-latitude ionosphere is more directly modified by  
182 magnetospheric processes that are largely controlled by the interaction between the solar

183 wind and the Earth's magnetosphere. In fact, many ionospheric phenomena that occur  
184 at high latitudes are footprint signatures of this interaction.

185 During the quiet periods, the different dynamics of the magnetic field fluctuations seem  
186 to describe physically different areas in the polar regions. The white profile, which delimits  
187 the spatial transition from an anti-persistent to a persistent dynamics, roughly corresponds  
188 to the boundary of the northern magnetic auroral oval. The spatial extension of this region  
189 changes with the geomagnetic activity level. In the bottom map of Figure 3, the local  
190 Hurst exponent values are reported for the disturbed geomagnetic activity level. The  
191 strong asymmetry in the daily character of the magnetic field's spatial fluctuations is still  
192 evident, but the region of anti-persistent character becomes smaller ( $05 < MLT < 17$ )  
193 than in the previous case, and the position of the white profile shifts towards geomagnetic  
194 latitude values less than  $70^\circ$  N in average, while the persistent region shows an increasing  
195 equatorward extent in the dark sector.

## 5. Summary and Conclusions

196 The time series of the geomagnetic field's horizontal component recorded by the Swarm  
197 A satellite were statistically analyzed in terms of scaling (fractal) features. The local  
198 Hurst exponent of the horizontal component was used to describe the spatio-temporal  
199 persistence character of the magnetic field fluctuations at high latitudes in the northern  
200 hemisphere during quiet and disturbed geomagnetic activity levels.

201 Maps of the local Hurst exponent give us the opportunity to localize the different lat-  
202 itudinal structures caused by different physical processes. The regions characterized by  
203 a larger value of the  $H$  exponent display magnetic field fluctuations with a larger spatio-

temporal coherence than the regions with a smaller  $H$ . Assuming that the observed fluctuations concern spatial fluctuations (i.e., under Taylor's hypothesis of frozen-in advected structures), because for a scale-invariant (fractal) signal there is a direct relationship between the Hurst exponent  $H$  and the exponent  $\beta$  of the power spectral density ( $S(k) \sim k^{-\beta}$  where  $\beta = 2H + 1$  and  $k$  is the wavenumber), we can suggest that the ionospheric polar regions characterized by different  $H$  values have different behaviors of the spectral density  $S(k)$  with the wavenumber  $k$ . Consequently, the geomagnetic field's spectral density is characterized by  $\beta > 2$  ( $\beta \sim 2.5 - 2.6$ ) in those regions where the Hurst exponent assumes values greater than  $\frac{1}{2}$  (for instance in the polar cap and auroral oval), while the spectral exponent is  $\beta < 2$  ( $\beta \sim 1.6 - 1.8$ ) in those regions with  $H$  values smaller than  $\frac{1}{2}$  (such as in the dayside of high and mid latitude regions).

This interpretation is partially supported by similar spectral regimes in the long-wavelength domain observed analyzing the turbulence properties of the electric field fluctuations at various altitudes over the auroral zone and polar cap [Weimer *et al.*, 1985; Heppner *et al.*, 1993; Golovchanskaya *et al.*, 2006; Kozelov *et al.*, 2008] where substantial differences between the scaling features of the fluctuations are not found in the two regions, as in our case [Golovchanskaya *et al.*, 2010]. However, there is not a simple correlation between the electric and magnetic fields, so that the results obtained for the electric field cannot be straight forwardly considered valid for the magnetic field. Sugiura *et al.* [1982] demonstrated a remarkably good correlation between the north-south component of the electric field ( $E_x$ ) and the east-west component of the magnetic field ( $B_y$ ) and Weimer *et al.* [1985] revealed that the Fourier spectra of these two components exhibited similar

226 power laws. Anyway, there is not a general one-to-one correspondence between the elec-  
 227 tric field  $E_x$  and magnetic one  $B_y$ , being this correspondence valid only in the auroral  
 228 oval, where the magnetic field lines are closed [*Smiddy et al.*, 1980; *Golovchanskaya et al.*,  
 229 2006]. Moreover, the present analysis treats the horizontal component of the geomagnetic  
 230 field and not of the single  $B_y$  component, and it describes an average result obtained using  
 231 about 3000 crossings without diversifying into different interplanetary magnetic field con-  
 232 ditions. Thus, our findings seem to be consistent with those by other authors suggesting  
 233 that the scaling properties of the magnetic field fluctuations in the auroral oval and in  
 234 the polar cap are basically the same. In contrast to earlier findings, however, our results  
 235 suggest that in these regions magnetic field fluctuations are characterized by power law  
 236 spectra with exponents  $\beta > 2$  at all the scales analyzed. Indeed, *Golovchanskaya et al.*  
 237 [2006, 2010] found that the power spectra of the electric field exhibit power-law relations  
 238 characterized by a break at scale  $\sim 40$  km, such that the slope  $\beta$  of the spectrum is greater  
 239 ( $\beta > 2$ ) for scales smaller than 40 km and lower ( $\beta < 2$ ) for scales over 40 km.

240 The different character of the spectral features that we find analyzing the magnetic field  
 241 fluctuations in the high-latitude polar regions may be a consequence of distinct turbulent  
 242 regimes: *i*) a strong shear-flow turbulence regime with  $\beta \sim 3$  (as in the polar cap and  
 243 auroral oval); *ii*) a shear-flow turbulence with an inverse energy cascade and/or a strong  
 244 gradient drift or current convective turbulence, which are both characterized by a spectral  
 245 index  $\beta \sim 5/3$  (as in the dayside high-latitude regions). In this picture, because for the  
 246 disturbed activity level we observe an expansion of the region characterized by a Hurst  
 247 exponent  $H > 1/2$ , this could be a consequence of plasma transport increase from distant

248 equatorial magnetotail regions during magnetic substorms, causing an expansion of the  
249 region affected by strong shear flow turbulence due to an enhancement of plasma convec-  
250 tion. We remark that the emergence of turbulence must be related to the enhancement  
251 of plasma density irregularities and Kelvin-Helmholtz instability.

252 In conclusion, this study shows the potential of ESA's Swarm mission to reveal the  
253 physical character of ionospheric high-latitude turbulence and its dependence on the ge-  
254 omagnetic activity level.

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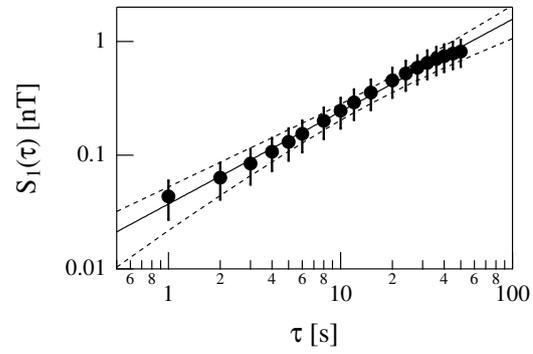
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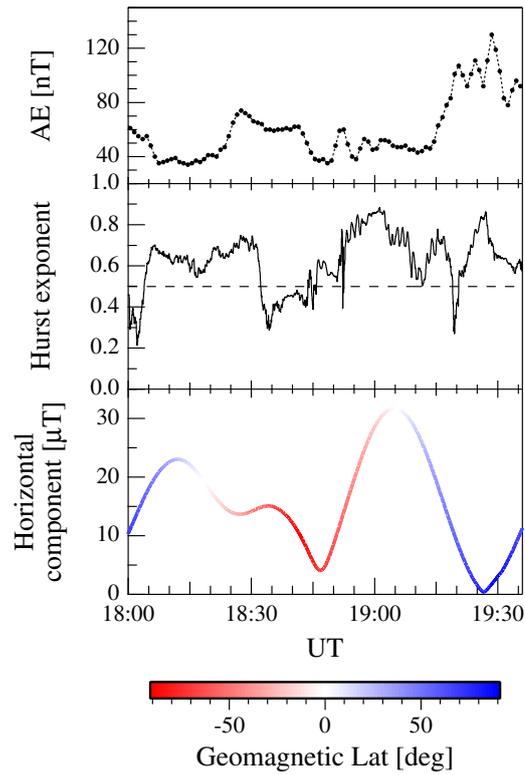
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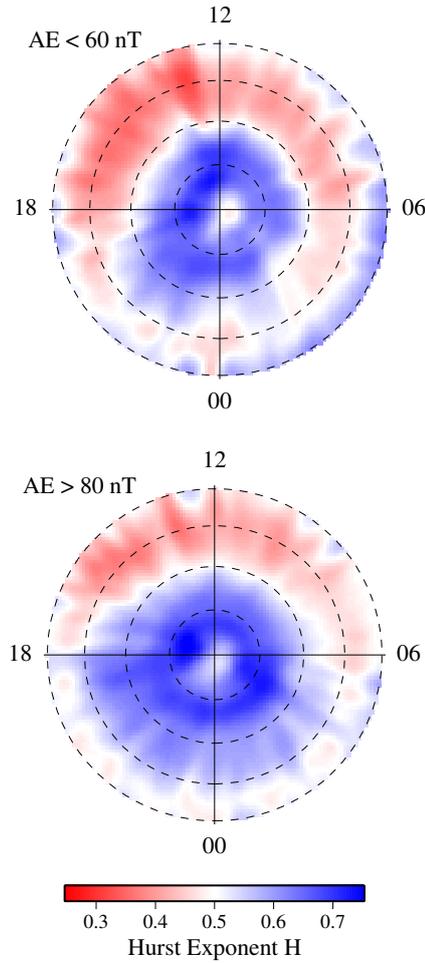
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**Figure 1.** A sample of the behavior of the 1<sup>st</sup> order structure function  $S_1(\tau)$ . The solid line is a power-law fit. Dashed lines show the 95% confidence interval.



**Figure 2.** From top to bottom: the Auroral Electrojet (AE) index, the local Hurst exponent ( $H$ ) and the magnetic field horizontal component for an orbit of the Swarm A satellite on April 1<sup>st</sup>, 2014. Dashed horizontal line (intermediate panel) is for  $H = 0.5$  and the color scale reported in the bottom plot shows the geomagnetic latitude of Swarm A during the selected orbit.



**Figure 3.** The local average Hurst exponent values in the northern hemisphere in a polar representation of magnetic local time (MLT) and geomagnetic latitude. Top and bottom maps refer to quiet ( $AE < 60$  nT) and disturbed ( $AE > 80$  nT) periods, respectively. Dashed circles are located at geomagnetic latitudes of  $50^\circ$  N,  $60^\circ$  N,  $70^\circ$  N and  $80^\circ$  N.