



# Intermittency and Passive Scalar Nature of Electron Density Fluctuations in the High-Latitude Ionosphere at Swarm Altitude

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## Key Points:

- The ionospheric electron density fluctuations show an intermittent and passive scalar behavior
- Plasma irregularities exhibit a complex dynamic behavior similar to a turbulent velocity field
- Turbulence is a key point to understand the occurrence of large local electron density gradients

## Plain Language Summary

Turbulence phenomena represent a key element for understanding the ionospheric dynamics and have a great impact on the features of both magnetic and electric fields and plasma density. It has been suggested that turbulence can be one of the mechanisms responsible for the generation of ionospheric irregularities, which may strongly affect the satellite navigation, positioning and communication systems. In such a framework the characterization of turbulent fluctuations of ionospheric plasma parameters paves the way for a better modeling of space weather related phenomena.

The ESA-Swarm mission provides a unique opportunity to investigate the relevance of turbulence phenomena in the topside ionosphere. Using data recorded on board Swarm A satellite at high latitudes, the intermittent and passive scalar features of electron density fluctuations are investigated. The obtained results support the idea that plasma density irregularities exhibit a complex dynamic behavior similar to that of a turbulent field. Turbulence is therefore confirmed as a fundamental element to be taken into account in order to fully understand the occurrence of large local electron density gradients.

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This article has been accepted for publication and undergone full peer review but has not been through the copyediting, typesetting, pagination and proofreading process which may lead to differences between this version and the Version of Record. Please cite this article as doi: 10.1029/2020GL089111

## Abstract

Some physical processes due to Sun-Earth interaction can influence the configuration and the dynamics of the high-latitude ionospheric plasma, particularly during geomagnetically disturbed periods. A possible consequence of this interaction is the occurrence of turbulent fluctuations that can be observed both in magnetic and electric fields and plasma density. Here, we focus on the intermittent nature of high-latitude ionospheric electron density fluctuations during geomagnetically disturbed periods as observed by one of the satellites of the ESA-Swarm constellation. The most obvious finding emerging from this study is the strong intermittent character of electron density fluctuations and the existence of an agreement between the anomalous scaling features of electron density fluctuations and those expected from a passive scalar quantity in fluid turbulence. This latter result supports the view of a passive scalar behavior of electron density in the high-latitude ionosphere that can have significant implications in the field of Space Weather studies.

## 1 Introduction

A large and growing body of literature has established that most space plasma processes involving interplanetary, magnetospheric and ionospheric plasma media, exhibit a turbulent and intermittent character (*Bruno and Carbone, 2013; Carbone et al., 2018; De Michelis et al., 2015; Kintner and Seyler, 1985; Tsunoda, 1988*). Turbulence plays a central role in several physical processes occurring in the interplanetary and near-Earth space regions and affects the plasma transport from the interplanetary space to the Earth's magnetosphere and ionosphere. Concerning the Earth's ionosphere, the turbulent nature of ionospheric plasma can play an important role in the formation and dynamics of plasma irregularities (*Hysell, 2020*). These can have a great impact on those technologies primarily related to telecommunications and satellite-based positioning systems. Indeed, plasma density irregularities can strongly influence the quality of electromagnetic signals that propagate in space. They are expected to be responsible for the ionospheric plasma scintillations, which strongly affect the Global Navigation Satellite System (GNSS), e.g., the Global Positioning System (GPS), and radio propagation (*Kelley, 1990; Knepp and Coleman, 2020; Wheelon, 1960*). Thus, it is of great importance to study and provide a better understanding of the processes that can affect ionospheric magnetized plasma density. In the past, it was clearly shown how in the high-latitude polar ionosphere the electrostatic turbulent fluctuations can significantly influence the plasma features, and in particular the electron density generating irregularities in the ionospheric medium (*Kintner and Seyler, 1985; Pécseli, 2015*). As clearly stated by *Kintner and Seyler (1985)*, the ionospheric electron density (and more in general the plasma density) may respond to the generation of electrostatic turbulence as a passive scalar quantity. In such a framework, the study of the features of electron density fluctuations can provide information on the physical processes responsible for the rise of turbulence in the ionospheric plasma. Several previous studies suggested different mechanisms for the rise of turbulence in the high-latitude ionosphere, such as, for instance, the gradient drift instability (GDI), the convective current instability (CCI), the Kelvin-Helmholtz instability (KHI), etc. (see, e.g., *Basu et al., 1990; Carlson et al., 2007; Cerisier et al., 1985; Keskinen et al., 1988; Mounir et al., 1991*). Thus, the investigation of the scaling features and intermittency of electron density fluctuations in the ionosphere are of central interest to unveil the origin of the plasma irregularities. The characterization of these features during geomagnetically disturbed periods is the central issue of this work.

The statistical properties of fluctuations in turbulent media (fluids and/or plasmas) can be investigated via the classical approach of scaling analysis of structure functions (see, e.g., *Frisch, 1995; Schmitt and Huang, 2016*), which are particularly suitable in studying scaling features dealing with turbulent signals and in investigating the occurrence of intermittency (anomalous scaling features and multifractality).

83 Given a time series  $\theta(t)$ , the  $q$ -th order generalized structure function at scale  $\tau$ ,  
84  $S_q(\tau)$ , is defined as

$$S_q(\tau) = \langle |\theta(t + \tau) - \theta(t)|^q \rangle, \quad (1)$$

85 where the angular brackets  $\langle \dots \rangle$  correspond with an ensemble average. In the case  
86 of scale-invariant signals/media the  $S_q$  behaves like a power law, being  $S_q(\tau) \propto \tau^{\zeta(q)}$   
87 where  $\zeta(q)$  is the corresponding scaling exponent. This is, for instance, what can be ob-  
88 served analyzing fluctuations in fluid and magnetohydrodynamic turbulence in the in-  
89 ertial range. The scaling properties of observed fluctuations and the set of scaling ex-  
90 ponents  $\zeta(q)$  allow characterizing the turbulence and quantifying the degree of intermit-  
91 tency, offering the possibility to discern among all the different models available to re-  
92 produce the observed scalings (*Frisch, 1995; Warhaft, 2000*). For example, according to  
93 the Kolmogorov's theory (K41) of homogeneous and isotropic turbulence, the scaling ex-  
94 ponent associated with the first-order structure function is expected to be  $\zeta(1) = 1/3$   
95 in the inertial range and the scaling exponents associated with higher-order structure  
96 functions are expected to scale linearly with the moment order  $q$  (*Frisch, 1995; Schmitt*  
97 *and Huang, 2016*) according to the following relation:

$$\zeta(q) = q\zeta(1) = \frac{q}{3}. \quad (2)$$

98 When Eq. (2) holds, one scaling exponent is enough to characterize the scaling features  
99 of fluctuations at all orders, and the employed K41-like models describe a mono-fractal  
100 behavior. However, real observations often reveal a departure from the linear dependence  
101 of scaling exponents  $\zeta(q)$  on the moment order  $q$  (*Frisch, 1995; Schmitt and Huang, 2016*).  
102 In this case, the system is characterized by anomalous scaling (multifractal features), which  
103 is the evidence of an intermittency phenomenon (*Frisch, 1995*). Of course, the higher  
104 the deviation from the straight line in Eq. (2), the stronger the degree of intermittency  
105 of the system.

106 The study of transport and mixing properties of vector and scalar fields, passively  
107 advected by the turbulent flow, is of great interest in fluid turbulence. For instance, the  
108 behavior of temperature and density fluctuations in a turbulent flow has been extensively  
109 investigated (*Falkovich et al., 2001; Falkovich and Sreenivasan, 2006; Shraiman and Sig-*  
110 *gia, 2000; Warhaft, 2000*) and modeled via the advection-diffusion equation,

$$\frac{\partial \theta}{\partial t} + (\mathbf{v} \cdot \nabla)\theta = \chi \nabla^2 \theta + f(\mathbf{r}, t), \quad (3)$$

111 where  $\theta$  is the passive scalar quantity,  $\mathbf{v}(\mathbf{r}, t)$  is the velocity field which transports the  
112 passive scalar quantity,  $\chi$  is the diffusivity and  $f$  is a forcing term necessary to attain  
113 a stationary state. According to this equation the passive scalar is a quantity driven by  
114 the flow and such to not affect the flow dynamics. Despite Eq. (3) is linear, complex fea-  
115 tures may be generated due to the coupling between  $\theta$  and  $\mathbf{v}$ . In particular, in the case  
116 of passive quantities transported by turbulent flows it has been shown that, conversely  
117 to the case of zero molecular diffusion which implies a rigid advection of the scalar field,  
118 a small amount of diffusivity gives rise to a different behavior (*Sreenivasan, 1991*). In-  
119 deed, the observed behavior is similar to what happens in a turbulent flow in the in-  
120 ertial range, where scalar fluctuations cascade from large to small scales at a constant scale-  
121 independent rate  $\epsilon_\phi$ .

122 Although the Kolmogorov-Obukhov-Corrsin (KOC) theory (*Corrsin, 1951; Obukhov,*  
123 *1949*) for passive scalar advection/diffusion predicts a simple scaling for the scalar in-  
124 crement structure functions  $S_q(\delta r) = (S_2(\delta r))^{q/2}$  (where  $\delta r$  is a specific spatial scale),  
125 deviations from this simple scaling can be observed. They are generally interpreted as

126 due to intermittency effects. The studies of the scaling features of passive scalar quan-  
 127 tities (such as, for instance, the temperature) in turbulent fluid flows (see e.g. Ref. *An-*  
 128 *tonia et al. (1984)*; *Ruiz-Chavarria et al. (1996)*) evidenced that the anomalous scaling  
 129 of a passive quantity increments is more pronounced than that of the velocity field. This  
 130 strong intermittent character of passive quantities is generally discussed in terms of a  
 131 double intermittency correction for kinetic energy and passive quantity fluctuations dis-  
 132 sipation rates (*Ruiz-Chavarria et al., 1996*).

133 Here, the scaling features of electron density and its passive scalar nature are in-  
 134 vestigated applying the structure function analysis on the ionospheric electron density  
 135 measured at high latitudes by one of the three satellites of the ESA-Swarm constella-  
 136 tion during two geomagnetically disturbed periods. In particular, a comparison between  
 137 the scaling features of electron density and those of passive scalar quantities in turbu-  
 138 lent media is presented. A brief discussion on the relevance of this study for the magnetosphere-  
 139 ionosphere coupling and its impact on the field of Space Weather studies are also pro-  
 140 vided.

## 141 2 Data set

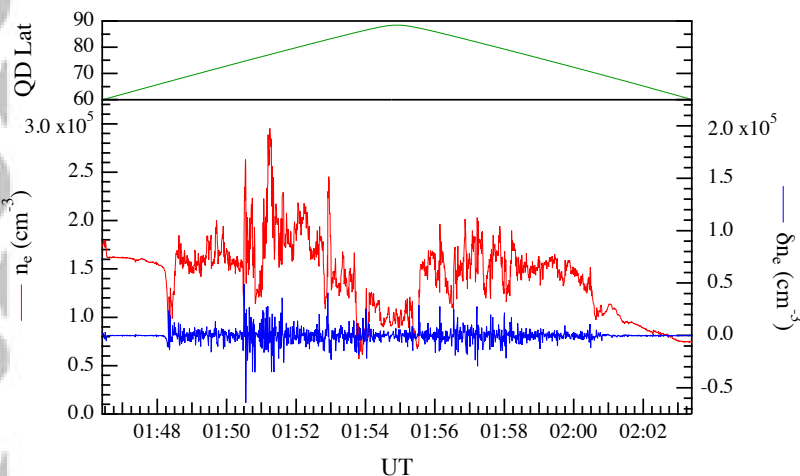
142 To investigate the scaling features and the intermittent nature of the high-latitude  
 143 electron density fluctuations we use the plasma measurements provided by the Langmuir  
 144 probes on-board Swarm A satellite. This satellite, which is one of the three satellites of  
 145 Swarm constellation, flies in a near-polar orbit at an altitude of about 460 km allowing  
 146 the study of those processes that mainly occur in the topside F region of the ionosphere.  
 147 We consider local electron density,  $n_e$ , measurements collected during the high-latitude  
 148 crossings of Northern ionosphere (QD Lat  $> 60^\circ$  N) in two different geomagnetically dis-  
 149 turbed days (17 March 2015 and 25 October 2016). Each crossing lasts approx. 20 min  
 150 (i.e., approx 2k points) and considering that Swarm A completes 15 orbits in a day, the  
 151 measurement dataset consists of  $\sim 60k$  points. Here, we use the quasi-dipole (QD) mag-  
 152 netic reference system (*Laundal and Richmond, 2017*). This latitude range has been cho-  
 153 sen to focus on the auroral oval and polar cap regions where plasma precipitation is ex-  
 154 pected to greatly increase during geomagnetically disturbed periods. The two selected  
 155 days refer to periods of mid-high geomagnetic disturbance, being the root-mean squared  
 156 value of the auroral electrojet geomagnetic index (AE) (*Davis and Sugiura, 1966*) for 17  
 157 March 2015 and 25 October 2016 equal to  $AE_{RMS} \simeq 830$  nT and  $AE_{RMS} \simeq 750$  nT,  
 158 respectively. Electron density measurements at a rate of 2 Hz have been selected from  
 159 the ESA ftp repository (<ftp://swarm-diss.eo.esa.int>).

160 Figure 1 shows a sample of the electron density,  $n_e$ , measurements and the corre-  
 161 sponding increments  $\delta n_e$  (where  $\delta n_e(\delta t) = n_e(t + \delta t) - n_e(t)$ ) calculated for  $\delta t = 0.5$  s  
 162 relative to one of the high-latitude crossings of the polar ionosphere on 17 March 2015.  
 163 Similar behaviors are observed for the other considered crossings. A clear increment of  
 164 electron density fluctuation amplitudes can be observed over the high latitudes.

## 165 3 Analysis and Results

166 We start our study analyzing the power spectral features of electron density fluc-  
 167 tuations.

168 Figure 2 displays the average Power Spectral Densities (PSDs) of the electron den-  
 169 sity  $n_e$  measured during the high-latitude crossings of the northern ionosphere in the two  
 170 selected periods. Each PSD is normalized by dividing it for the variance  $\sigma_{n_e}^2$  of the elec-  
 171 tron density  $n_e$  of the correspondent period. Both PSDs are characterized by a power-  
 172 law behavior,  $S(f) \sim f^{-\alpha}$ , with a scaling exponent  $\alpha \sim 5/3$  over more than two or-  
 173 ders of magnitude. These spectral features can be considered the signature of a turbu-  
 174 lent character of the observed fluctuations, and agree with previous findings (see, e.g.:



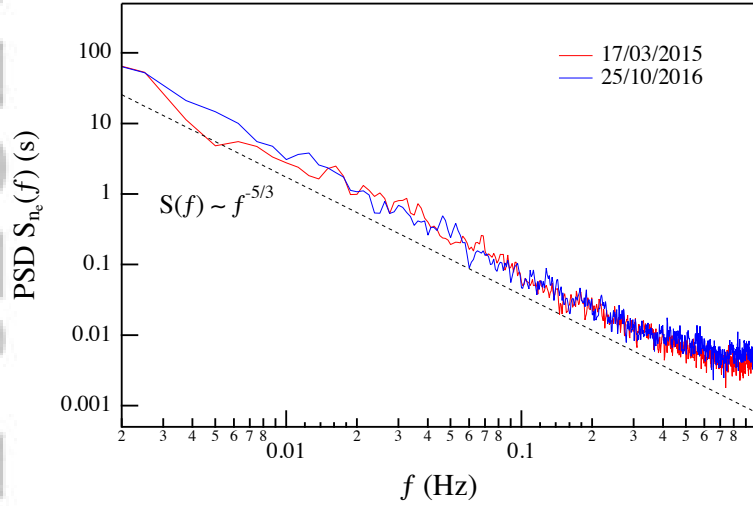
**Figure 1.** A sample of electron density,  $n_e$ , and electron density increments  $\delta n_e$  at the scale  $\delta t = 0.5$  s during one of the high-latitude crossings of the Northern ionosphere on 17 March 2015. The upper panel shows the quasi-dipole magnetic latitude (QD Lat).

175 *Basu et al.*, 1990; *Kelley et al.*, 1980; *Kintner and Seyler*, 1985; *Spicher et al.*, 2015, and  
 176 references therein) and theoretical predictions (*Kintner and Seyler*, 1985). Indeed, al-  
 177 though the analysis is done in the temporal domain, we can assume that the observed  
 178 fluctuations are mainly spatial being the low-frequency temporal fluctuations principally  
 179 the effect of Doppler-shifted and stationary spatial variations (*Basu et al.*, 1990; *Con-*  
 180 *solini et al.*, 2020). Under this assumption the observed PSDs are compatible with what  
 181 is expected in the case of fluid strong turbulence at least for wavenumbers  $k = f/v_s$   
 182 in the range  $[3, 60] \cdot 10^{-3}$  km, being the satellite velocity  $v_s \sim 7.8$  km/s.

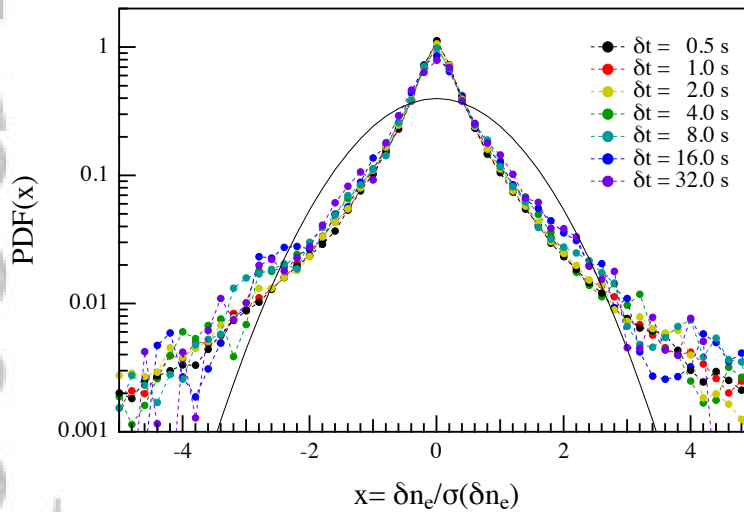
183 Considering that the PSD spectral exponents are in the range  $\alpha \in (1, 2)$ , the sig-  
 184 nal is clearly nonstationary but its increments are stationary (weak stationarity) (*Con-*  
 185 *solini et al.*, 2020; *Davis and Sugiura*, 1966) and consequently the requirements of struc-  
 186 ture function analysis are satisfied.

187 Another interesting feature of electron density fluctuations/increments,  $\delta n_e(\delta t)$ , is  
 188 their non-Gaussian statistics at short timescales ( $\delta t < 50$  s). Figure 3 shows the Prob-  
 189 ability Density Functions (PDFs) of the electron density increments normalized to the  
 190 standard deviation ( $\sigma_{\delta n_e}$ ) for timescales  $\delta t < 50$  s as obtained considering the whole  
 191 dataset. In particular, the normalization of the increments to the standard deviation is  
 192 necessary to remove the dependence on the different geomagnetic activity level. The PDFs  
 193 exhibit a large departure of  $\delta n_e$  from the Gaussian statistics, being the PDFs charac-  
 194 terized by a leptokurtic shape (i.e., a distribution with a kurtosis higher than 3) (*De-*  
 195 *Carlo*, 1997). The PDFs of the electron density increments are analogous to those found  
 196 by evaluating the PDFs of velocity increments in turbulent fluid flows in the low-end in-  
 197 ertial range (*Frisch*, 1995), where large departures from Gaussian statistics are observed  
 198 at small scales towards the dissipation scale. Furthermore, the PDFs display a shape that  
 199 is dependent on the timescale  $\delta t$ , i.e., the shape of the PDF is not scale invariant (PDFs  
 200 collapsing is poor). This feature is a signature of intermittency, as it occurs in turbu-  
 201 lence. We will return on this point later in the discussion.

202 We first investigate the occurrence of scaling features in structure functions by study-  
 203 ing their dependence on the moment order  $q$ . The left panel of Figure 4 displays the de-  
 204 pendence of  $q^{th}$ -order structure functions of electron density on time increments  $\delta t$ . We  
 205 limit the investigation of scaling features to moment orders  $q \leq 4$  due to the available



**Figure 2.** The average Power Spectral Density of the electron density  $n_e$  for the two selected periods. The dashed line is a power law with a scaling exponent  $-5/3$ , plotted for comparison.



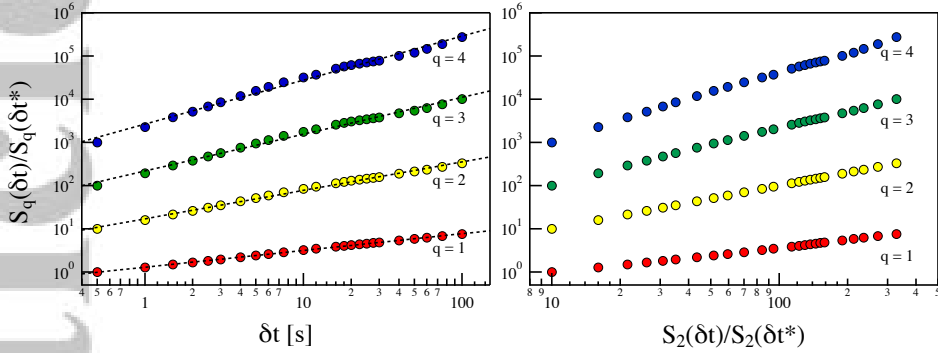
**Figure 3.** The PDFs of the electron density increments,  $\delta n_e$ , for  $0.5 \text{ s} \leq \delta t \leq 32 \text{ s}$ . Data have been scaled to the respective standard deviation. The black curve is a unit-variance Gaussian distribution for reference.

number of measurements (see, e.g., *Dudok de Wit, 2004*). A clear power-law scaling,  $S_q(\delta t) \sim \delta t^{\zeta(q)}$  is found over more than 2 orders of magnitude, supporting that electron density fluctuations/increments have a self-affine structure over a wide range of scales, and thus they show scale-invariance.

Let us now move to the relative scaling of the  $q^{\text{th}}$ -order structure function,  $S_q(\delta t)$ , versus the  $2^{\text{nd}}$ -order one,  $S_2(\delta t)$ , which is expected to scale as follows:

$$S_q(\delta t) = (S_2(\delta t))^{\xi(q)}, \quad (4)$$





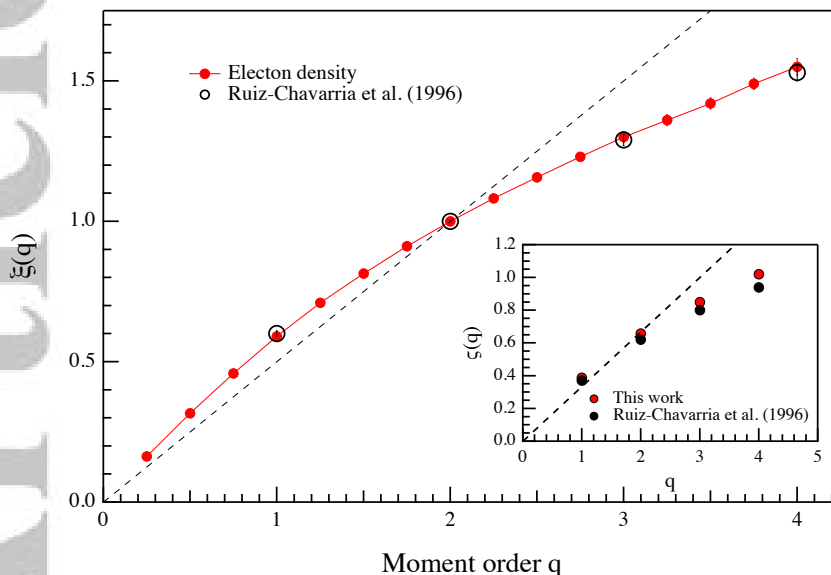
**Figure 4.** Left Panel: Scaling of structure function,  $S_q(\delta t)$ , as a function of  $\delta t$  for  $q = 1, 2, 3$ , and 4. Dashed lines are nonlinear power-law best fits. Right Panel: The relative scaling of the  $q^{\text{th}}$ -order structure function,  $S_q(\delta t)$ , as a function of the  $2^{\text{nd}}$ -order one,  $S_2(\delta t)$ . Structure functions are normalized to the corresponding value for timescale  $\delta t^* = 0.5$  s and scaled by a factor  $10^{q-1}$  for convenience.

212 and according to *Benzi et al.* (1993) is named Extended Self-Similarity (ESS) analysis.  
 213 For mono-scaling (monofractal) signals  $\xi(q)$  is expected to scale linearly with the moment  
 214 order  $q$ , being  $\xi(q) = \zeta(q)/\zeta(2) = q/2$ . Differently, anomalous scaling features,  
 215 i.e., a non-linear dependence on  $q$ , are expected in the case of multi-scaling (multifractal)  
 216 signals, i.e., signals whose scaling features are no-longer homogeneous but acquire  
 217 a local dependence with a spreading of local scaling exponents (*Benzi et al.*, 1984; *Schmitt*  
 218 *and Huang*, 2016).

219 The right panel of Figure 4 displays the relative scaling of the  $q^{\text{th}}$ -order structure  
 220 function,  $S_q(\delta t)$  of the electron density increments  $\delta n_e(\delta t)$  with  $\delta t \in [0.5, 100]$  s ver-  
 221 sus the  $2^{\text{nd}}$ -order one,  $S_2(\delta t)$ . The observed dependence of  $q^{\text{th}}$ -order structure  
 222 functions on the  $2^{\text{nd}}$ -order one agrees with the predictions provided by Eq. 4 over 1.5 decades.

223 To characterize the relative scaling features we evaluate the relative scaling expo-  
 224 nents  $\xi(q)$  for the  $q^{\text{th}}$  order electron density structure functions versus the  $2^{\text{nd}}$ -order one.  
 225 The best fit is a power-law dependence between  $S_q(\delta t)$  and  $S_2(\delta t)$  for timescales in the  
 226 interval  $\delta t \in [0.5, 50.0]$  s, which correspond to spatial scales in the range  $[4, 400]$  km,  
 227 taking into account the satellite speed and assuming the fluctuations to be mainly spa-  
 228 tial (see the previous discussion on PSD). Figure 5 shows the behavior of the relative scal-  
 229 ing exponents,  $\xi(q)$ , as a function of the moment order  $q$ . We considered also non-integer  
 230  $qs$  in order to better trace the scaling exponents trend. A clear nonlinear dependence  
 231 of these exponents is found, suggesting that we are in presence of anomalous scaling fea-  
 232 tures of the structure functions. This is the signature of the occurrence of intermittency  
 233 in the case of electron density fluctuations.

234 In order to unveil if the observed anomalous scaling can have the same character  
 235 of that observed in the case of passive scalar quantities in fluid turbulence, we have over-  
 236 plotted in Figure 5 the relative scaling exponents computed using results on scaling fea-  
 237 tures of temperature fluctuations in fluid turbulence, which has been proven to be a pas-  
 238 sive scalar quantity. The values of the temperature scaling exponents come from the work  
 239 by *Ruiz-Chavarría et al.* (1996), and they are listed in their Table 2. The agreement is  
 240 excellent. Conversely, the observed anomalous scaling features seem to be different from  
 241 those observed in *Spicher et al.* (2015). This could be due to the different range of scales,  
 242 where the anomalous scaling features are investigated, being the analysis in *Spicher et*  
 243 *al.* (2015) done at smaller scales,  $\delta t \leq 0.1$  s.



**Figure 5.** The relative scaling exponents  $\xi(q)$  of the  $q^{\text{th}}$ -order structure function,  $S_q(\delta t)$ , of the electron density measurements. Empty circles refer to the same quantities computed using values from *Ruiz-Chavarria et al. (1996)*. The dashed line shows the expected behavior in the case of absence of intermittency. The inset shows the comparison among the observed scaling exponents  $\zeta(q)$  of the  $q^{\text{th}}$ -order structure function for this work and those reported in *Ruiz-Chavarria et al. (1996)*; dashed line is Kolmogorov's K41 scaling,  $\zeta(q) = q/3$ .

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To better underline the agreement between our results and those from *Ruiz-Chavarria et al. (1996)*, we report in Table 1 the relative scaling exponents  $\xi(q)$  of electron density fluctuations and the corresponding values computed from *Ruiz-Chavarria et al. (1996)*. The agreement between the relative scaling exponents is very striking, suggesting that the intermittency observed in the case of electron density fluctuations is of the same universality class of intermittency observed in the case of a passive scalar quantity in fluid turbulence. We remark that up to the 4<sup>th</sup> moment order the measured values of the scaling exponents,  $\zeta(q)$ , display a quasi-universal character as clearly shown in *Warhaft (2000)* (see Figure 11 therein) where data from different experiments are compared.

$q$	$\xi(q)$ This work	$\xi(q)$ from <i>Ruiz-Chavarria et al. (1996)</i>
1	$0.59 \pm 0.01$	$0.60 \pm 0.01$
2	1	1
3	$1.30 \pm 0.02$	$1.29 \pm 0.02$
4	$1.55 \pm 0.03$	$1.53 \pm 0.03$

**Table 1.** The relative scaling exponents,  $\xi(q)$ , of electron density fluctuations as obtained in the current work, compared to those computed from data in *Ruiz-Chavarria et al. (1996)* for temperature fluctuations. The value of the 2<sup>nd</sup> order structure function scaling exponent in our case is  $\zeta(2) = [0.658 \pm 0.008]$ .



## 4 Conclusions

The main findings of our analysis are: 1) the signature of the occurrence of intermittency in electron density fluctuations in the auroral oval ionospheric region, 2) electron density fluctuations have the same universality class of a passive scalar quantity in fluid turbulence. To our knowledge this is the first observational evidence of the passive scalar nature of electron density in high-latitude ionospheric regions.

Turbulence has been claimed to be a very relevant phenomenon in the ionospheric medium. It can strongly affect the plasma density and, in particular, the electron density generating strong irregularities that may have a great impact on satellite navigation, positioning and communication systems. Several works (see, e.g.: *Kintner and Seyler, 1985; Pécseli, 2015*) have stressed the role that fluid-like turbulence generated by  $\mathbf{E} \times \mathbf{B}$  gradient drift, current convective, shear flow and/or Kelvin-Helmholtz instabilities can play in generating plasma irregularities. Nevertheless it is necessary to take into consideration that also the intermittency of passive scalar quantities in fluid turbulence can be very relevant and generally stronger than that observed in the velocity field (see, e.g.: *Kraichnan, 1994; Ruiz-Chavarria et al., 1996; Warhaft, 2000*).

Thus, our results on the intermittency and the passive scalar nature of the ionospheric electron density fluctuations strongly support the idea that turbulence is probably the most relevant phenomenon capable of generating plasma irregularities, i.e., multiscale patchy structures responsible for the occurrence of ionospheric scintillations and radio propagation anomalies in the ionospheric medium. The observed universality character of electron density scaling features and passive scalars in fluid turbulence also provides an indication that plasma velocity field is turbulent in the high-latitude ionosphere. Moreover, the great similarity between the observed scaling exponents of the electron density structure functions and those of passive scalars in fluid turbulence suggests that the turbulent nature of plasma velocity field may be fluid. This point supports the idea that this turbulent field may be generated by strong  $\mathbf{E} \times \mathbf{B}$  drift velocity shears/gradients.

Further studies are clearly necessary in order to put our findings in a general theory for the formation of ionospheric plasma irregularities.

## Acknowledgments

The results presented rely on data collected by ESA-Swarm mission. We thank the European Space Agency that supports the Swarm mission. Swarm electron density measurements can be accessed at the ESA ftp repository (<ftp://swarm-diss.eo.esa.int> (see also <http://earth.esa.int/swarm>). The authors kindly acknowledge V. Papitashvili and J. King at the NSSDC of the Goddard Space Flight Center for the use permission of 1 min OMNI data and the NASA CDAWeb team for making geomagnetic indices data available at <https://cdaweb.gsfc.nasa.gov/index.html/>. The authors acknowledge financial support from European Space Agency (ESA contract N. 4000125663/18/I-NB-“EO Science for Society Permanently Open Call for Proposals EOEP-5 BLOCK4” (INTENS)) and from the Italian MIUR-PRIN grant 2017APKP7T on “Circumterrestrial Environment: Impact of Sun-Earth Interaction”.

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