

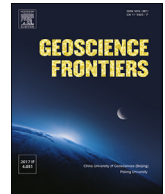
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Research Paper

## Scattering and absorption imaging of a highly fractured fluid-filled seismogenic volume in a region of slow deformation

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## ABSTRACT

Regions of slow strain often produce swarm-like sequences, characterized by the lack of a clear mainshock-aftershock pattern. The comprehension of their underlying physical mechanisms is challenging and still debated. We used seismic recordings from the last Pollino swarm (2010–2014) and nearby to separate and map seismic scattering (from P peak-delays) and absorption (from late-time coda-wave attenuation) at different frequencies in the Pollino range and surroundings. High-scattering and high-absorption anomalies are markers of a fluid-filled fracture volume extending from SE to NW (1.5–6 Hz) across the range. With increasing frequency, these anomalies approximately cover the area where the strongest earthquakes occurred from the sixteenth century until 1998. In our interpretation, the NW fracture propagation ends where carbonates of the Lucanian Apennines begin, as marked by a high-scattering and low-absorption area. At the highest frequency (12 Hz) the anomalies widen southward in the middle of the range, consistently marking the faults active during the recent Pollino swarm. Our results suggest that fracture healing has closed small-scale fractures across the SE faults that were active in the past centuries, and that the propagation of fluids may have played a crucial role in triggering the 2010–2014 Pollino swarm. Assuming that the fluid propagation ended at the carbonates barrier in the NW direction, fractures opened new paths to the South, favoring the nucleation of the last Pollino swarm. Indeed, the recently active faults in the middle of the seismogenic volume are marked by a high-scattering and high-absorption footprints. Our work provides evidence that attenuation parameters may track shape and dynamics of fluid-filled fracture networks in fault areas.

## 1. Introduction

The southern Apennines and the Calabrian arc are among the most seismically-active areas of the Mediterranean region. Despite the intense seismic activity, gaps in historical documentation (Scionti et al., 2006) and the low population density of the area (Cinti et al., 1997) have marked the Pollino range as a "seismic gap" between the southern Apennines and the Crati Valley (Rovida et al., 2011; Tertulliani and Cucci, 2014) (Fig. 1). Paleoseismological investigations show clear evidence of large-magnitude earthquakes ( $M_{6.5-7}$ ) occurred along two seismogenic areas, the Pollino and Castrovallari faults (POL and CAS in Fig. 1) (Cinti

et al., 1997, 2002). A few earthquakes of moderate magnitude occurred during seismic sequences struck these areas in the past centuries, as the 1693 Pollino seismic sequence (Tertulliani and Cucci, 2014) or the 1708 earthquake ( $M_{5.5}$ ). In 1998, a  $M_L = 5.0$  earthquake occurred in the Mercure basin, northwest of the Pollino massif, triggering the Pollino-Mercure basin sequence (Guerra et al., 2005). Although the uncertainties on the location of historical events are higher than for recent seismicity, there is evidence of a migration of the seismogenic volumes from SE to NW across the Pollino range between 1559 and 1998 (Fig. 1). Between 2010 and 2014 a sequence of almost 10,000 earthquakes of small-to-moderate magnitude with two mainshocks of  $M_L 4.3$  and  $M_L 5.0$

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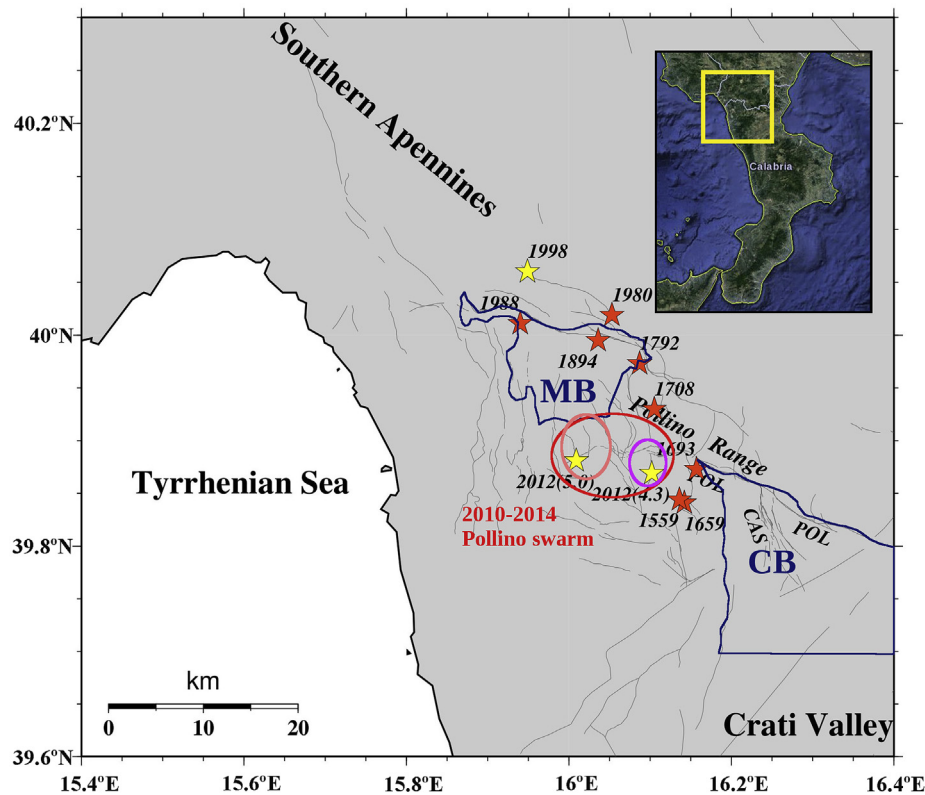
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**Fig. 1.** Map of the Pollino transition zone, located between the southern Apennines and the Crati Valley (yellow rectangle in the top right panel). Orange stars represent the presumed epicenters of historical earthquakes until 1988, yellow stars represent mainshocks occurred in the area in the last 20 years. MB and CB, bordered by blue lines, are the Mercuri Basin and the castrovillari Basin, respectively. CAS and POL show the Castrovillari and Pollino Faults. The red ellipse surrounds the area of the 2010–2014 seismic swarm, which occurred in two slightly separated areas (smaller ellipses). Thin gray lines are the detailed faults from Brozzetti et al. (2017) in the Pollino area and from ITHACA catalogue for the outermost faults surrounding the Pollino area.

occurred in the Pollino area (Passarelli et al., 2015; Totaro et al., 2015) (Figs. 1 and 2). During this sequence, characterized by a migration in time from NW toward SE (Fig. 1), a slow-slip event lasting several months (between 2012 and 2013) was detected in the same area via the GPS monitoring network (Cheloni et al., 2017). Seismic data have been used in combination with geological and remote-sensing data to map seismically-active normal faults (Brozzetti et al., 2017) and to evaluate local site effects (Napolitano et al., 2018) in the area. The relationship between slow-strain areas, often characterized by a combination of seismic brittle failures and aseismic slip, and swarm-like seismic activity is one of the possible scenarios proposed, e.g., by Lohman and McGuire (2007), Peng and Gombert (2010), and Passarelli et al. (2015). Even though the correlation between swarm-like activity and their fault mechanisms is not fully understood, these works suggest that aseismic processes are a general and common feature driving swarm-like sequences. For this reason, novel geophysical imaging is necessary to improve hazard estimation for earthquakes that, in slow-strain regions, are episodic and spatially migrating (Landgraf et al., 2017).

Seismic attenuation tomography has the potential to image the extension of highly-fractured volumes, especially when fluids saturate them inducing seismicity (De Siena et al., 2016, 2017; Amoroso et al., 2017). The two physical mechanisms that induce inelastic attenuation while a wave travels through the crust are seismic scattering and absorption. The introduction of the radiative transfer equation solved by numerical Monte Carlo simulations put forward the important role of multiple scattering in the generation of coda waves (Paasschens, 1997; Sato et al., 2012). In their seminal paper, Gusev and Abubakirov (1999) devised a strategy to invert the broadening envelopes (the seismic intensities recorded at a station, Fig. 3) for transport turbidity, a single parameter describing scattering attenuation and back-scattering of coda waves. At crustal scale, the peak-delay time (i.e. the time lag from the direct-wave onset to the maximum amplitude of the signal envelope) is generally considered a direct measure of multiple forward scattering after correcting for path-propagation (Saito et al., 2002; Takahashi et al., 2007; Calvet and Margerin, 2013).

Coda waves (the later portion of the seismogram) are the main manifestation of the redistribution of seismic energy caused by multiple scattering. Their attenuation ( $Q_c^{-1}$ ) is measured from the exponential decay of coda energy envelopes with time (Aki and Chouet, 1975). At long lapse times and in a uniform anisotropic half-space, coda waves theoretically enter the diffusive regime, which in turn implies equality between coda attenuation and absorption (Shapiro et al., 2000). This assumption is valid at regional scale for seismicity constrained in the crust, i.e. in a thick layer regime (Calvet et al., 2013; Margerin, 2017) – which in turn means that earthquakes are mostly constrained inside a thick ( $\leq 30$  km) crust. At late lapse times, coda waves from these earthquakes have undergone multiple scattering interactions: equipartition between P- and S-waves (Hennino et al., 2001) is considered a good marker of the diffusion regime in this setting. Calvet and Margerin (2013) demonstrated that coda-wave attenuation at  $t_w = 80$  s is a measurement of seismic absorption for epicentral distances between 0 and 90 km in the Pyrenees. Borleanu et al. (2017) used a similar assumption to image Vrancea (Romania), revealing extension and shape of sedimentary basins. The joint use of peak-delays and  $Q_c^{-1}$  in a thick crust like the Pollino area (average Moho depth at 45 km) thus allows to separate scattering attenuation from absorption.

$Q_c^{-1}$  regionalisation can reconstruct the geotectonic characteristics of seismogenic regions (e.g., Ugalde et al., 2002). Nevertheless, while peak-delays are sensitive to a tight area around the seismic rays (Saito et al., 2002), assigning  $Q_c^{-1}$  to ray paths between source and receiver is inaccurate (Del Pezzo et al., 2016). Also, the regionalisation approach hinders testing of anomalies as it does not use a proper forward model for inversion. Recently, modelling of coda amplitude through kernel functions has been proposed (Margerin et al., 2015; Del Pezzo et al., 2016) and tested (Mayor et al., 2016; De Siena et al., 2017) at crustal and local (volcanic) scales. Most of these functions are computed via a Monte Carlo numerical simulation of the Energy Transport Equation. With the above-mentioned assumptions, the multiple anisotropic scattering process leads to diffusion in a time-window where equipartition takes place

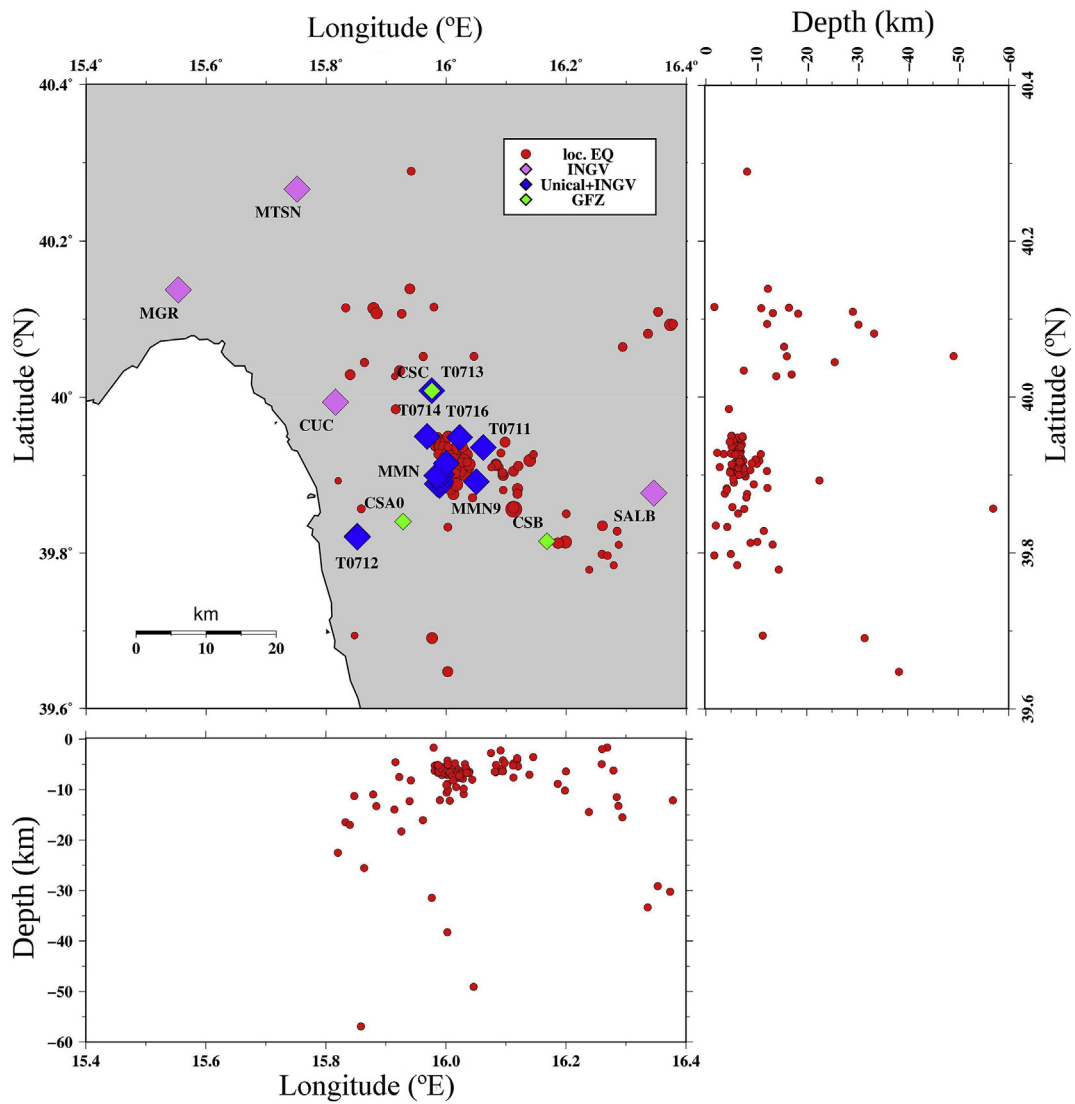


Fig. 2. Map of the study area showing earthquakes used in this work (red circles) and seismic stations divided in three groups: permanent and temporary stations operated by the *Istituto Nazionale di Geofisica e Vulcanologia, INGV* (purple; Margheriti et al., 2013); temporary stations installed by GFZ (green, FDSN network code 4A; Passarelli et al., 2012); permanent and temporary stations operated by *Università della Calabria* (blue).

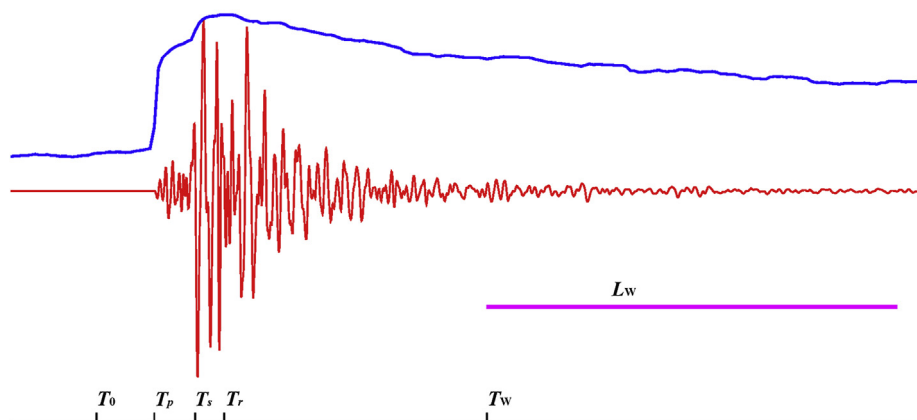


Fig. 3. Example of a waveform (red line) used in this work and its envelope (blue line).  $T_p$  and  $T_s$  are the P- and S-wave arrival times,  $T_r$  is the time at which the waveform reaches the maximum value,  $T_w$  is the beginning of the analysis window and  $L_w$  is its length. All these times are referenced to the origin time  $T_0$ .

(Hennino et al., 2001; Souriau et al., 2011) and in the absence of leakage and strong boundary conditions (De Siena et al., 2013; Margerin, 2017). All studies agree that in this case the functions have two maxima at source and station positions and expand around them depending on the average scattering properties of the medium (e.g., Obermann et al., 2013; Del Pezzo et al., 2016).

We used earthquakes of the 2010–2014 Pollino sequence and crustal earthquakes recorded in the surrounding area (lon: 15.4°E, 16.4°E; lat: 39.6°N, 40.4°N) to produce 2D maps showing the spatial variations of peak-delay time and  $Q_c^{-1}$ . While we regionalized peak-delay measurements, we used an analytic approximation of the diffusive sensitivity kernels (Del Pezzo et al., 2016) to model the effective amplitude decrease of coda waves at late lapse times, and inverted  $Q_c^{-1}$  in space (De Siena et al., 2017). We tested the resolution and reliability of the final results by inverting on grids having different steps, changing damping parameters and performing checkerboard tests. We use these maps as a proxy of the lateral variations of seismic scattering and absorption in the Pollino area, allowing us to image fault structures, fracture networks, and their influence on fluid propagation and seismicity between the 16th century and today.

## 2. Geological and geophysical settings

The Pollino area (southern Italy) is a transition zone between the Southern Apennines NE-verging collision and the Calabrian rollback subduction zone. Located on the northern side of the Calabrian fore-arc and accretionary wedge, it is a striking example of the faster subduction (relative to the normal subduction in the Mediterranean) of the African under the European plate (Faccenna et al., 1996). The Calabrian fore-arc collided with the continental margins of Nubia, forming the Maghrebides, and of Apulia, forming the Apennines. Two shallow tectonic units coexist in the Pollino transition zone: (1) the allochthonous *Liguride Unit*, representing the remnants of the northern continental margin of Neotethys (Cello et al., 1996) and (2) the *Apennines Platform*, a thick carbonate shelf succession represented by the Verbicaro Unit overlaying the Pollino Unit. Furthermore, in the Campotenese area, the superposition of the Verbicaro unit onto the Campotenese-Pollino unit is marked by a ductile shear zone along which the Jurassic–lower Miocene carbonate sequence of the Pollino unit is strongly deformed.

In this complex geological setting, medium to strong earthquakes ( $4 \leq M \leq 6$ ) occur on the northwest flank of the Pollino Range and Mercure basin area on roughly SE–NW trending normal faults (Brozzetti et al., 2017). The earthquakes accommodate the arc-normal northeast–southwest extension. Left lateral strike-slip regional faults oriented WNW–ESE (Van Dijk et al., 2000) and alignments of East- and West-dipping normal faults have been found in the area. Meanwhile, on the Tyrrhenian coastal side, regional East-dipping normal faults have been recognized across north Calabria.

Fig. 1 shows a map of a wide area including the Pollino Range and its surroundings. Here, we show the known faults taken from both the detailed mapping of Brozzetti et al. (2017) for the Pollino range, and ITHACA catalogue for the outermost main structures, as well as other geological features considered in the interpretation. The same figure shows the epicenters of monitored (yellow) and historical (orange) earthquakes occurred in the area (taken from Ferranti et al., 2017), and the epicentral area of the last seismic sequence (2010–2014, red ellipse). The locations of historical earthquakes are achieved using the Boxer method (Gasparini et al., 2010) based on the evaluation of the epicenter as the center of mass of largest intensities. Even if this methodology was tested using more recent epicenters, the uncertainties that afflict these locations are much higher than those of monitored events, in the order of some km. Nevertheless, taken into account the location errors, epicenters follow an SE–NW trend through time, at least until the 1998 Mercure basin earthquake.

## 3. Data and methods

We have analyzed velocity waveforms data recorded by permanent and temporary seismic stations operating in the Pollino area during the 2010–2014 seismic sequence. We selected 117 crustal earthquakes of local magnitude ( $M_L$ ) ranging from 1.8 to 4.3, source-receiver distance between 1 km and 72 km, and depths between 2 km and 56 km. We performed the waveform selection to obtain a satisfactory coverage of crossing rays across the area. However, while hundreds of events were available in the central part, most of them being located roughly in the same spot, fewer were found in other volumes around the 2010–2014 sequence. Most of the selected events ( $\sim 80\%$ ) are located at depth smaller than 10 km (in Supplementary Materials Fig. S1). To increase the ray coverage filling gaps were left by shallower we used deeper crustal earthquakes; nevertheless, only 7 of them ( $< 2\%$ ) are deeper than 30 km.

Since the seismic sequence continued for years at a variable rate, the number of temporary stations often changed in the area. The selected data set includes 21 seismic stations, some of which were permanent while others were operated from months to some years by three institutions: (1) *Università della Calabria*; (2) *Istituto Nazionale di Geofisica e Vulcanologia (INGV)*; (3) *GFZ* (Fig. 2). Data of INGV and GFZ stations were downloaded from the EIDA online database (Bianchi et al., 2015). The distribution of seismic stations and earthquakes is crucial to test the reliability of our results. Since we used data coming from three different databases, the selected earthquakes were located before any other analysis.

We selected 911 waveforms with clear P-wave phase and coda-to-noise ratio higher than 3 at 30 s from the origin time in the frequency range of 1–32 Hz. The final dataset also comprises 357 waveforms with clear picking of both P- and S-waves, the first used for pick-delay mapping. We bandpass-filtered the whole seismograms in 5 frequency bands (1–2 Hz, 2–4 Hz, 4–8 Hz, 8–16 Hz, 16–32 Hz) applying a Butterworth filter of order 4, forward and backwards. We computed envelopes from the absolute value of the Hilbert transform of the signal. We then smoothed them with a moving window of length 8 times the inverse central frequency – an average between what is used between regional (Calvet and Margerin, 2013) and local volcanic (De Siena et al., 2016) scales.

### 3.1. Peak-delay time measurement and mapping

We measured the peak-delay time as the lag between the P-wave onset and the maximum S-wave amplitude (Takahashi et al., 2007). The base- $\log_{10}$  of the peak-delay time ( $t_r^T(f)$ , in seconds) is related to the base- $\log_{10}$  of the epicentral distance ( $R$ , in km) in each frequency band  $f$  by:

$$\log_{10}(t_r^T(f)) = A_r(f) + B_r(f)\log_{10}(R) \quad (1)$$

where  $A_r$  and  $B_r$  are the coefficients of the fit (in Supplementary Material Fig. S2). The variations of the peak-delay time with respect to these trends ( $\Delta\log_{10}(t_r(f))$ ):

$$\Delta\log_{10}(t_r(f)) = \log_{10}(t_r(f)) - \log_{10}(t_r^T(f)) \quad (2)$$

do not depend on geometrical propagation. They are assumed as a measure of the relative strength of accumulated S wave scattering along each ray path (Saito et al., 2002). In other words, high values of  $\Delta\log_{10}(t_r(f))$  mean that the ray path crosses regions of high heterogeneity (Calvet et al., 2013). For the mapping, we selected the area shown in Figs. 1 and 2 (lon: 15.4°E, 16.4°E; lat: 39.6°N, 40.4°N) and divided it into rectangular blocks of size  $0.05^\circ \times 0.05^\circ$ . We thus evaluated peak-delays along 2D source-receiver rays and assigned their values to each node with a regionalisation approach (Takahashi et al., 2007) (Fig. 4). Even though no exact forward model exists, we tested the stability of the results by repeating the same analysis in a rectangular grid of steps  $0.1^\circ \times 0.1^\circ$

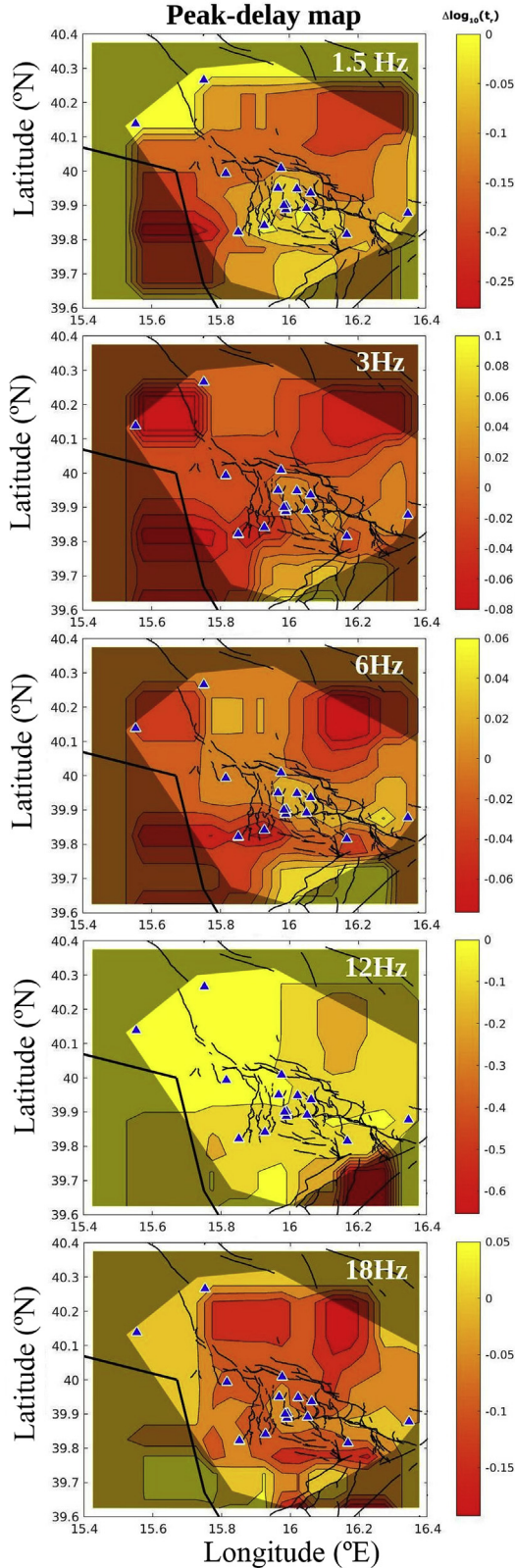


Fig. 4. Peak-delay map. Central frequency is shown on the top-right for each panel, blue triangles represent seismic stations used in this work. Yellow colors represent high-scattering zones, red colors represent low-scattering zones.

(Fig. S2). We then interpreted only patterns recovered by the test and where  $Q_c^{-1}$  is recovered by resolution tests (see next section).

### 3.2. Inverse coda-Q measurement and mapping

We consider the energy envelope decay of coda waves described by Aki and Chouet (1975) using the following equation:

$$E(t, f) = S(f)t^{-\alpha} \exp\left(\frac{-2\pi ft}{Q_c}\right) \quad (3)$$

where  $E(t, f)$  is the power spectral density,  $S(f)$  is a source/site term,  $t$  is the lapse-time from the origin of the event, and  $Q_c^{-1}(f)$  is the frequency-dependent inverse coda quality factor.  $\alpha$  is equal to 3/2 in a layer characterized by an anisotropic multiple scattering regime (Paasschens, 1997; Calvet et al., 2013). The choice of the coda window is crucial to map stable lateral variations of seismic attenuation. We want to exclude the increasing transient regime occurring at short lapse time as shown by Calvet and Margerin (2013) while we require a signal-to-noise ratio greater than 5 across the whole window of analysis. Coda windows in our work start at lapse time  $t_w = 15$  s from the origin-time of the selected earthquake and are characterized by a length of the time window  $L_w = 10$  s (Fig. 3). This coda window provides envelopes with sufficient energy for measuring  $Q_c^{-1}$  at four of the five previously-mentioned frequency bands (Fig. 5). At 18 Hz, a significant number of envelopes present a signal-to-noise ratio lower than 5; we considered the results at this frequency as unreliable and analyzed only the remaining frequencies.

In our time window, equipartition is achieved after 15 s (blue line in Fig. 5) while  $Q_c^{-1}$  shows no consistent variations with epicentral distance (in Supplementary Materials Fig. S3). In the absence of leakage (Calvet and Margerin, 2013; Margerin, 2017) we can thus interpret  $Q_c^{-1}$  as a measurement of absorption. Still, we need to acknowledge the effect of leakage on our data. Our scale is similar to that of Gaebler et al. (2015) ( $0.5^\circ \times 0.5^\circ$ ), who estimate a transport mean free path of  $l = 50\text{--}110$  km between 3 Hz and 24 Hz, approximately our frequency range. The ratio between the transport mean free path and the average crustal thickness  $H = 45$  km (Piana Agostinetti and Amato, 2009) thus varies between  $H = 0.4$  km and  $H = 1$  km. The effect of leakage is thus not negligible but still well above the limit ( $H = 0.2$  km) where equipartition brakes down (Margerin, 2017).

We used a non-linear approach to solve Eq. (3) in the other four frequency bands (1–2 Hz, 2–4 Hz, 4–8 Hz, 8–16 Hz). We divided the envelope of seismic coda waves between 15 s and 25 s in windows of length 2 s for each source-station pair. Then, we normalized the energy in each smaller window dividing by that of the last window to remove the source effect,  $S(f)$ . We used a grid search on 1000 trial  $Q_c$  values to find the best solution. In each window, we subtracted the normalized energy and the model, then minimized the residual L1 norm for all time windows. This non-linear solution is more stable than the standard linearized technique for coda signals having low signal-to-noise ratios (Ibanez et al., 1993).

We mapped the sensitivity of the  $Q_c^{-1}$  parameters to space  $(x, y)$  by computing diffusive sensitivity kernels (Del Pezzo et al., 2016) (in Supplementary Material Fig. S4). For each source–receiver pair of coordinates  $(x_s, y_s)$  and  $(x_r, y_r)$ , the kernels are defined as follows:

$$f[x, y, x_r, y_r, x_s, y_s] = \frac{1}{4\pi\delta_x D^2 \delta_y} \exp\left[-\frac{(x - \frac{x_r + x_s}{2})^2}{2(\delta_x D)^2} + \frac{(y - \frac{y_r + y_s}{2})^2}{0.5(\delta_y D)^2}\right] + \frac{1}{2\pi\delta_x D^2 \delta_y} \exp\left[-\frac{(x - x_s)^2}{2(\delta_x D)^2} + \frac{(y - y_s)^2}{2(\delta_y D)^2}\right] + \frac{1}{2\pi\delta_x D^2 \delta_y} \exp\left[-\frac{(x - x_r)^2}{2(\delta_x D)^2} + \frac{(y - y_r)^2}{2(\delta_y D)^2}\right] \quad (4)$$

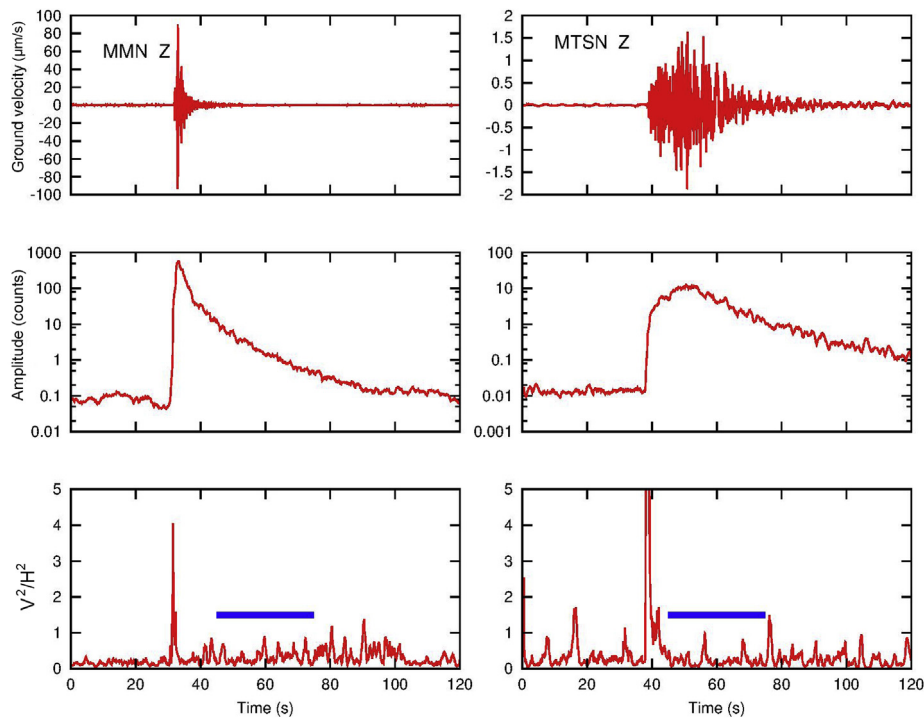


Fig. 5. Equipartition of energy shown for two earthquakes. The two earthquakes are characterized by small (4 km, on the left) and long (44 km, on the right) epicentral distance. From top to bottom: chosen seismograms; plot of the smoothed envelope and the mean noise amplitude; ratio between the kinetic energy of the vertical and horizontal components,  $V^2/H^2$ . Blue segments represent the windows chosen to compute coda analysis.

where  $\delta_x$  and  $\delta_y$  are the spatial apertures of the weighting functions along the  $x$  and the  $y$  axis. These apertures can be set to 0.2 for a wide range of crustal scattering parameters in a diffusive regime (Del Pezzo et al., 2016). We set the same area and grid defined for peak-delay analysis, inverting for the spatial distribution of  $Q_c^{-1}(f)$  in space in the four remaining frequency bands (Fig. 6). The mapping is performed with the inversion scheme described by De Siena et al. (2017), with the hypothesis that the energy is entirely lost inside the predefined grid. We evaluated effective stability and resolution of the solution by selecting the best damping parameter between the size of the regularized solution and its fit to the given data (L-curves in Fig. S5). As for peak-delay analysis, we tested the stability of our results with a coarser node spacing (Fig. S6). Finally, we tested our resolution performing checkerboard tests with multiple grid node spacing ( $0.05^\circ \times 0.05^\circ$  grid in Fig. S7 and  $0.1^\circ \times 0.1^\circ$  grid in Fig. S8).

### 3.3. Parameter space variation and mapping

We plotted scattering vs. absorption measurements in their parameter space after removing the mean computed over all measurements (De Siena et al., 2016). This graphical method makes results clearer on map, relating the relative variation of  $Q_c^{-1}(f)$  (horizontal-axis, absorption) to  $\Delta \log_{10}(t_r(f))$  (vertical-axis, scattering) and setting a color for each possible combination of the space parameter divided in 4 quadrants (Fig. S9). In this case, we set: red for high scattering and high absorption (HS-HA), orange for low scattering and high absorption (LS-HA), light blue for high scattering and low absorption (HS-LA), green for low scattering and low absorption (LS-LA). We additionally set gray for values with a level of discrimination less than 1% of the maximum variations. We characterize each node of the grid with the same palette (Fig. 7, left column).

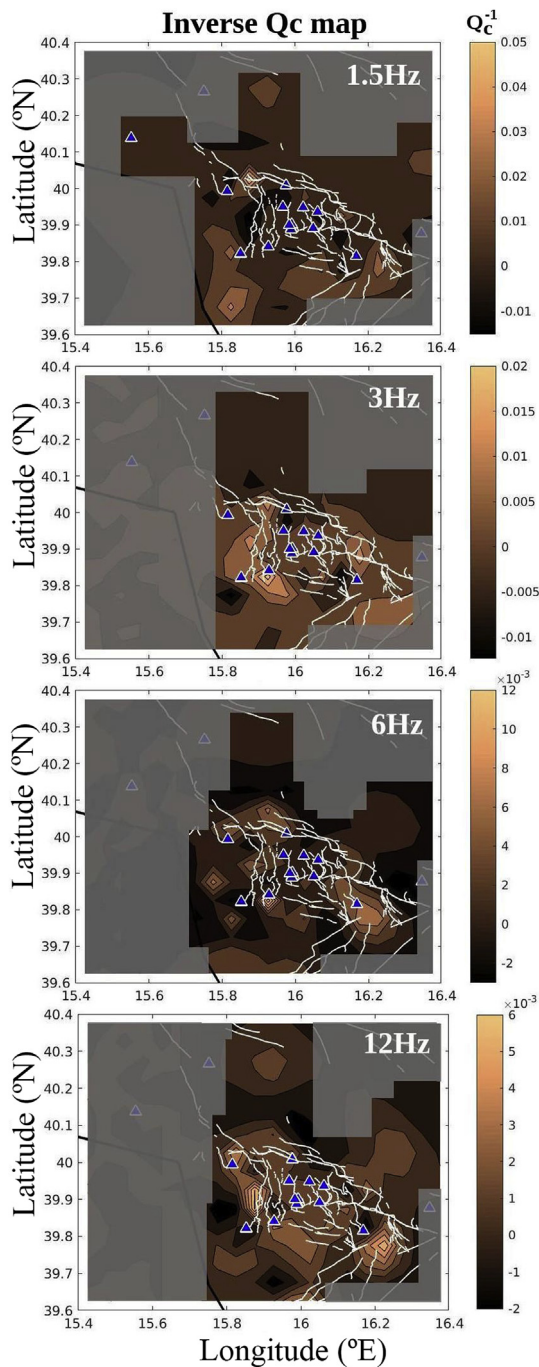
## 4. Results

Figs. 4 and S2 (in Supplementary Materials) show the regionalisation

results for peak-delays on 2D map characterized by grids with steps of  $0.05^\circ$  and  $0.1^\circ$ , respectively. Comparison between them shows the stability of patterns of dimension equal or greater than  $0.1^\circ \times 0.1^\circ$ . Fig. 6 and Fig. S6 (in Supplementary Materials) show the  $Q_c^{-1}(f)$  results obtained for the same grid steps. At low frequencies, late coda-to-noise ratios are higher thus providing results affected by lower uncertainties. In a given time window, the coda-to-noise ratio progressively decreases as the frequency increases, thus reducing data used for the analysis at higher frequencies and/or reliability of results. For each frequency band, we computed the L-curves (log-log plots of the norm of the regularized solution versus the norm of the corresponding residual norm). After testing different damping parameters, we chose 0.1 as the best compromise to obtain scattering and absorption maps at all frequencies.

### 4.1. Discussion

In Fig. 7 we present (left) and interpret (right) the results for each frequency band. Due to the depth of the Moho (Piana Agostinetti and Amato, 2009), we assume that multiple scattering acts in an anisotropic scattering layer. Most earthquakes are located at depths lower than 10 km, supporting this assumption. The effect of leakage may still be important from previous calculations (Margerin, 2017). Still, the Moho in the study area is mostly flat (Piana Agostinetti and Amato, 2009), minimizing the lateral variations produced by this mechanism. In summary, we can safely assume that our results are representative of the shallow crust and look for any relationships between the observed scattering (peak-delay) and absorption (coda attenuation) patterns and the geological features of the investigated area. We do not observe any orange zone (LS-HA) stable and big enough to be discussed. The Calabro-Lucania coastal range (western side of the maps—Fig. 7 left) changes from LS-LA at low frequency (green, 1.5 Hz and 6 Hz) to HS-LA (cyan) at 12 Hz. However, the resolution tests show that the area is unresolved. The same is true for one of the most stable patterns in frequency, i.e., the green LS-LA zone visible in the north-eastern sector of the maps.



**Fig. 6.** Inverse  $Q_c$  map. Central frequency is shown on the top-right for each panel, blue triangles represent seismic stations used in this work.

A cyan HS-LA pattern both stable in frequency and resolved by our mapping comprises the part of the Lucanian Apennines (named *LA* in Fig. 7 right) extending from Mount Sirino to Lake Pertusillo. Geologically, this is a compact carbonate stand-alone block, surrounded by faults and isolated through clay formations. Fractured volumes of minor importance are present in this block due to the high-compression regime in the region. This geological configuration likely explains the high scattering values at all frequencies. The low-absorption values are instead a consequence of the cohesion of the carbonate block, which acts as a shield between the extensive northern hydrocarbon deposit in the Val D'Agri zone and the southern Pollino area.

The red-coloured HS-HA patterns are the most remarkable and well-resolved of our study, and the most interesting for an interpretation in

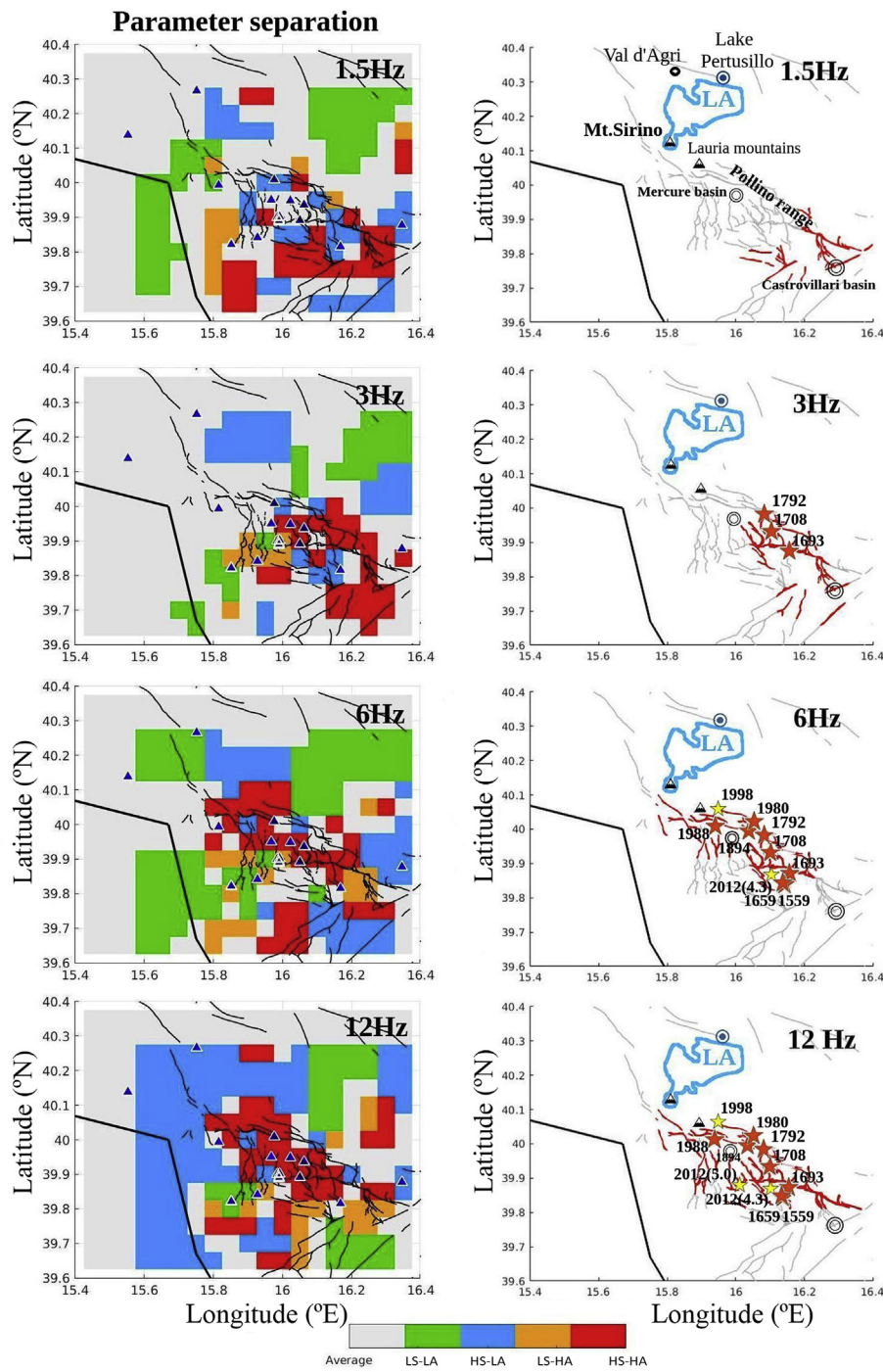
terms of fluid-filled fracture networks (Quintal et al., 2014). The faults located inside it are coloured red in the maps on the right side of Fig. 7. The Mercure and the Castrovillari basins (MB and CB, respectively) are both characterized by shallow-marine deposits. The former is filled with alluvial, fluvial-deltaic and lacustrine sediments (Ferranti et al., 2017); the latter is filled with clay, sand (in the shallower layer), fine marine sediments, and conglomerates. Both basins are likely characterized by high absorption, contributing to the red pattern. The HS-HA red pattern marks the fault system underneath the Pollino Range at 1.2 Hz. We suggest a relationship with the prominent geological characteristic of a multiple fragmented shallow-water carbonate succession. The Lauria mountains (LM) are the north-west extension of the carbonate succession of the Pollino Ridge, and they are characterized by the same HS-HA (red) at the same frequencies.

The HS-HA pattern changes its position with changing frequency. At 1.5 Hz, the red area marks the S-SE sector of the map, crossed by extensive fault networks; at 3 Hz the pattern is still present in the SE sector but becomes predominant in the Pollino area. The red color spreads across the Mercure basins and the Lauria mountains at 6 Hz, and fills the area of the 2010–2014 seismic swarm at 12 Hz. The relatively high-absorption values correlate well with the high  $v_p/v_s$  values found in the same area by Fig. 7 in Barberi et al. (2004), and by Totaro et al. (2015). The results confirm the scenario proposed for the 2010–2014 seismic swarm by Passarelli et al. (2015) of a large role played either by the increase of pore pressure or fluid infiltration within the seismogenic zone in favouring the aseismic slip by lowering the normal stress on the fault plane.

Historical earthquake locations are affected by significant uncertainties, differently from those occurred in the last 20 years. Even considering these uncertainties, the HS-HA pattern migrates with increasing frequency following the chronological order of historical earthquakes from the sixteenth century until Pollino swarm. The most recent earthquakes, e.g., the Mercure basin earthquake and the Pollino swarm itself, likely reactivate small fracture networks as fluids migrate along the main faults, modifying the way fluids permeate fault structures. Poro-elasticity studies proved the central role of fluid pressure diffusion between connected fractures, and showed how the behavior of the connected fracture system is frequency dependent (Quintal et al., 2014).

Hunziker et al. (2018) pointed out that the still-unknown relationship between seismic attenuation, diffusion of fluid pressure and fracture connectivity can be crucial to identify a highly-fractured volume and track its hydraulic behavior. Our results strongly suggest such a relationship in some of the fault areas. Changes of peak-delay and  $Q_c^{-1}$  with frequency mark in fact heterogeneity of different size (Saito et al., 2002): in the case of faults, small fractured volumes in active zones can only be imaged by shorter-wavelength (higher-frequency) attenuation parameters. While these fracture networks heal quickly, in seismically-active areas they continuously reactivate due to distant seismicity (Xue et al., 2009), at least where fluids are abundant. The healing of smaller fractures is instead permanent in areas not deformed by recent activity. Here, larger unhealed structures should be visible from their long-wavelength (low-frequency) high-scattering and high-absorption signatures. We thus infer that the migration of the HS-HA patterns with increasing frequency shows the effect of fracture healing (Xue et al., 2009).

The anomalies effectively track the migration of seismogenic zones since the sixteenth century (Fig. 1) by mapping scattering and absorption changes in frequency inside the fracture-network. While we already confirmed the role of fluids in nucleating historical and recent seismicity (Passarelli et al., 2015), our results support the view that the Pollino swarm represents a change in the usual seismicity trend of the area. The stop to the SE–NW migration of the HS-HA pattern at 12 Hz is, in our interpretation, due to the barrier represented by the high-scattering compressed carbonates under the Lucanian Apennines. Since the Pollino swarm, a permanent network has been installed in the area and may provide sufficient data to test our inferences in the presence of renewed seismic activity.



**Fig. 7.** Attenuation and scattering results (left) and schematic interpretation of the stable and resolved patterns (right) obtained in 4 frequency ranges. Palette used: red–high scattering/high absorption, orange–low scattering/high absorption, cyan–high scattering/low absorption and green–low scattering/low absorption. Blue triangles show the seismic stations. The cyan pattern corresponding to the Lucanian Apennines (shortly named LA) extends from the Mount Sirino (SW) to the Lake Pertusillo (NE). Faults are coloured and stars (orange for earthquakes before 1998, yellow for more recent earthquakes) are displayed each time they are located in a red high scattering/high absorption block.

## 5. Conclusions

Peak-delay and inverse coda quality factor have been measured together and mapped using a 2D approximation in four frequency bands to image the scattering attenuation and absorption properties of the slowly-deforming Pollino area (southern Italy). The pattern of anomalies obtained from our analysis shows a correlation with the geological features of the area and with the progressive shift in location of earthquakes. High-scattering and high-absorption anomalies (HS-HA) change their position from inactive faults in the south/southeastern Castrovillari basin (1.5 Hz) to the northwestern Mercure Basin and Pollino swarm area (6–12 Hz) across the fault systems (3 Hz). We infer that the Lucanian Apennine compressed carbonates, marked by high-scattering, constitute

a barrier for fracture network connectivity and fluid propagation. The two sedimentary basins are characterized by high absorption. However, high-absorption and high-scattering patterns observed at higher frequency likely reveal fluid-filled fractured volumes of shorter dimension.

At low frequency, the patterns map the structures of the fault systems inactive (1.5 Hz) or fractured by seismic sequences dating back to the 16<sup>th</sup>–18<sup>th</sup> centuries (3 Hz). Here, shorter fractures are progressively healed, with high scattering and absorption marking extended fault areas permanently deformed. At 6 Hz the anomalies characterize the NW branch of the faults, tracking the seismic activity of the last two centuries up to the Mercure basin earthquake (1998). The earthquake migration is likely due to the progressive increase of the contribution of small-scale fluid-filled fractures to scattering and absorption. At the highest



frequency (12 Hz) the HS-HA area becomes broader and covers the fault system that generated the last Pollino swarm (2010–2014), SW of the Mercure basin. We infer that the fracture propagation trend was blocked NW by the high-scattering Lucanian Apennine carbonates. The scattering and absorption picture described by our interpretation provides an independent view on the well-known geological features of the area, the known fault network, and the migration of seismogenic zones through the last five centuries, focusing on the important role of fluid saturation. Further observations with better distributed earthquakes and seismic stations may give deeper insight into the physics underlying these processes. More efforts are necessary to apply successfully the laboratory results to the Earth crust, that means in our case to measure or estimate the time scale of processes such as fracture healing and fluid migration throughout connected/non-connected fracture networks. Still, our work supports the view that attenuation and scattering parameters are able to reconstruct shape and dynamics of fluid-filled fracture networks in fault areas.

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### Appendix A. Supplementary data

Supplementary data to this article can be found online at <https://doi.org/10.1016/j.gsf.2019.09.014>.

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